

TU DARMSTADT

Electron Scattering as a Tool to Study Excited States with Unusual Features

- $J^{\pi} = 1/2^+$ resonance at neutron threshold in ⁹Be
- Structure of the Hoyle state in ¹²C

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Possible role of ⁹Be in the production of ¹²C



In *n*-rich environment (core-collapse supernovae) this reaction path may provide an alternative route for building up the heavy elements and triggering the *r* process

$J^{\pi} = \frac{1}{2}^{+}$ state at threshold



S-DALINAC



Lintott spectrometer



Detector system



- Si microstrip detector system: 4 modules, each 96 strips with pitch of 650 μm
- Count rate up to 100 kHz
- Energy resolution 1.5x10⁻⁴

Spectrum of the ⁹**Be(e,e**[']**) reaction and deconvolution**



⁹Be(γ,n) extracted from ⁹Be(e,e[^])



Comparison: ⁹Be(γ,n) and ⁹Be(e,e[´])



Reaction rate of $\alpha \alpha n \rightarrow {}^{9}Be$



Rate is shown as ratio to the compilation of C. Angulo et al., NPA 656, 3 (1999)

- For $T_9 = 0.1 3$ this resonance determines exclusively ${}^{4}\text{He}(\alpha,\gamma){}^{8}\text{Be}(n,\gamma){}^{9}\text{Be}$ chain
- Determined reaction rate differs up to 20% from adopted values

Form factor of the $J^{\pi} = \frac{1}{2}^{+}$ state



- NCSM (C. Forssén): correct q dependence but difference in magnitude compared to the data
- $B(C1) \neq B(E1)$ at photon point $k = q \rightarrow$ violation of Siegert theorem ?

Convergence behavior of NCSM results



Radial transition densities from NCSM



Problems and possible improvements

• E1 operator needs to be renormalized \rightarrow effective charges

but: convergence even slower than for the binding energies

• State is unbound \rightarrow put a tail onto the insufficient HO wave functions

↔ projection onto a (core + particle) subspace

successful for strong s.p. component and well defined core configuration, but both are questionable here

 Calculate also the transverse E1 response to investigate possible sources of the violation of Siegert's theorem

Hoyle state in ¹²C

- Astrophysical motivation
- High-resolution electron scattering experiments

• Structure of the Hoyle state: a "BEC" ?

Comparison with FMD and a-cluster model predictions

M. Chernykh, H. Feldmeier, T. Neff, PvNC and A. Richter, PRL 98, 032501 (2007)

Motivation: triple alpha process



Reaction rate:

$$\Gamma_{rad} = \Gamma_{\gamma} + \Gamma_{\pi} = \frac{\Gamma_{\gamma} + \Gamma_{\pi}}{\Gamma} \cdot \frac{\Gamma}{\Gamma_{\pi}} \cdot \Gamma_{\pi}$$

$$r_{3\alpha} \propto \Gamma_{rad} \exp\left(-\frac{Q_{3\alpha}}{kT}\right)$$

S.M. Austin, NPA 758, 375c (2005)

- Electron scattering $\rightarrow M(E0) \rightarrow extraction of \Gamma_{\pi}$
- 0⁺₁ and 0⁺₂ (7.654 MeV) states in ¹²C
 - density distributions
 - model predictions



Measured spectra



Antisymmetrized A-body state

$$|Q
angle\,=\,\mathcal{A}(|q_1
angle\otimes|q_2
angle\otimes\ldots\otimes|q_A
angle)$$

- Single-particle states

$$\langle \mathbf{x} | q
angle = \sum_{i} c_{i} \, \exp \Big[- rac{(\mathbf{x} - \mathbf{b}_{i})^{2}}{2a_{i}} \Big] \otimes |\chi_{i}^{\uparrow}, \chi_{i}^{\downarrow}
angle \otimes |\xi
angle$$

- Gaussian wave packets in phase space (a_i is width, complex parameter \mathbf{b}_i encodes mean position and mean momentum), spin is free, isospin is fixed
- Describes α -cluster states as well as shell-model–like configurations

UCOM interaction

- Derived form the realistic Argonne V18 interaction
- Adjusted to reproduce binding energies and charge radii of some "closed-shell" nuclei

α-cluster model

- FMD wave function restricted to α -cluster triangle configurations only
- "BEC" model
 - System of 3 ⁴He nuclei in 0s state (like α condensate)
 - Hoyle state is a "dilute gas" of α particles
- Volkov interaction
 - Simple central interaction
 - Parameters adjusted to reproduce α binding energy, radius, $\alpha-\alpha$ scattering data and ground state energy of ^{12}C
 - Only reasonable for ⁴He, ⁸Be and ¹²C nuclei

¹²C densities



Elastic form factor



- H. Crannell, data compilation
- Described well by FMD

Transition form factor to the Hoyle state



• H. Crannell, data compilation, plus measurements at low q from Darmstadt

- Described better by α -cluster models
- FMD might be improved by taking α - α scattering data into account

What is actual structure of the Hoyle state ?

Overlap with FMD basis states



- In the FMD and α-cluster model the leading components of the Hoyle state are cluster-like and resemble ⁸Be + ⁴He configurations
- But in the "BEC" model the relative positions of α clusters should be uncorrelated

Summary and outlook

Summary

- Hoyle state is not a true Bose-Einstein condensate
- ⁸Be + α structure
- Outlook
 - ¹²C: 0⁺₃ and 2⁺₂ states
 - ¹⁶O: 0+ states at 13 15 MeV
 - Improved E0 matrix element: Γ_{π} for decay of the Hoyle state

One-Level Approximation of *R***-Matrix Theory**

$$\sigma_{\gamma,n}(E_{\gamma}) = \frac{\pi}{2 k_{\gamma}^2} \frac{2J+1}{2I+1} \frac{\Gamma_{\gamma} \Gamma_n}{\left(E_{\gamma} - E_{\lambda} - \Delta(E_{\gamma})\right)^2 + \frac{\Gamma^2}{4}}$$

$$\Gamma_{\gamma} = \frac{16 \pi}{9} e^2 k_{\gamma}^3 B(E1, k) \downarrow \qquad \Gamma_n = 2 \sqrt{\epsilon (E_{\gamma} - S_n)}$$
$$\Delta(E_{\gamma}) = -\gamma^2 (S(E_{\gamma}) - B_n) = 0$$
$$\Gamma_n \gg \Gamma_{\gamma} \qquad \Gamma \simeq \Gamma_n$$

$$\sigma_{\gamma,n}(E_{\gamma}) = \frac{16\pi^2}{9} \frac{e^2}{\hbar c} g_J B(E1,k) \downarrow \frac{E_{\gamma} \sqrt{\epsilon (E_{\gamma} - S_n)}}{(E_{\gamma} - E_R)^2 + \epsilon (E_{\gamma} - S_n)}$$
$$\Gamma_R(E_R) = 2\sqrt{\epsilon (E_R - S_n)}$$

One-Level Approximation of *R***-Matrix Theory**

$$\sigma_{e,e'}(E_x) = const \times \frac{E_x \sqrt{\epsilon} (E_x - S_n)}{(E_x - E_R)^2 + \epsilon (E_x - S_n)}$$

$$\Gamma_R = 2\sqrt{\epsilon(E_R - S_n)}$$

Extraction of the B(C1,k) Transition Strength

$$\sqrt{B(C1,q)} = \sqrt{B(C1,0)} \left(1 - q^2 \frac{R_{tr}^2}{10} + q^4 \frac{R_{tr}^4}{280} - \ldots\right)$$

Siegert's theorem: at photon point (q = k)

$$B(E1,q) = (k/q)^2 B(C1,q)$$
$$B(E1,k) = B(C1,k)$$

Extraction of the B(E1) Transition Strength

