

Electromagnetic properties of nuclei: from few- to many-body systems

Lecture 11

Many-body methods - Applications

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November 23rd, 2017

Lecture series for SFB 1245
TU Darmstadt

S.B. *et al.*, PRL **111**, 122502 (2013)

We have developed a new many-body technique that allows to study break up observables from first principles and is built on the ground of a true many-body method, hence has the potential to surpass previous limits

$$(\bar{H} - z^*)|\tilde{\Psi}_R(z^*)\rangle = \bar{\Theta}|\Phi_0\rangle$$

The first implementation we provided is based on the CCSD approximation scheme

$$T = T_1 + T_2$$

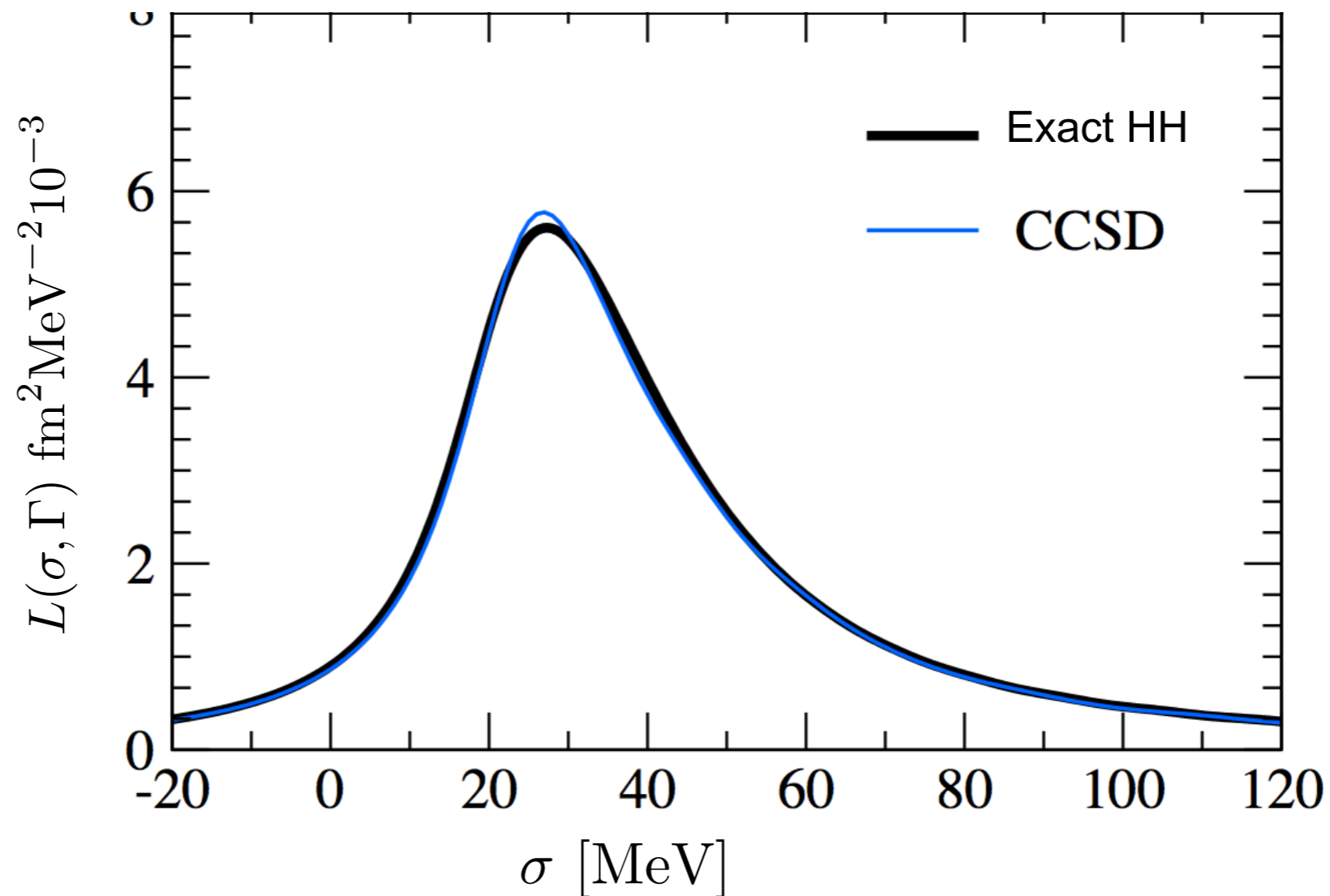
$$R = R_1 + R_2$$

S.B. *et al.*, Phys. Rev. Lett. **111**, 122502 (2013)

Lorentz Integral Transform with NN forces at N³LO

Validation ⁴He

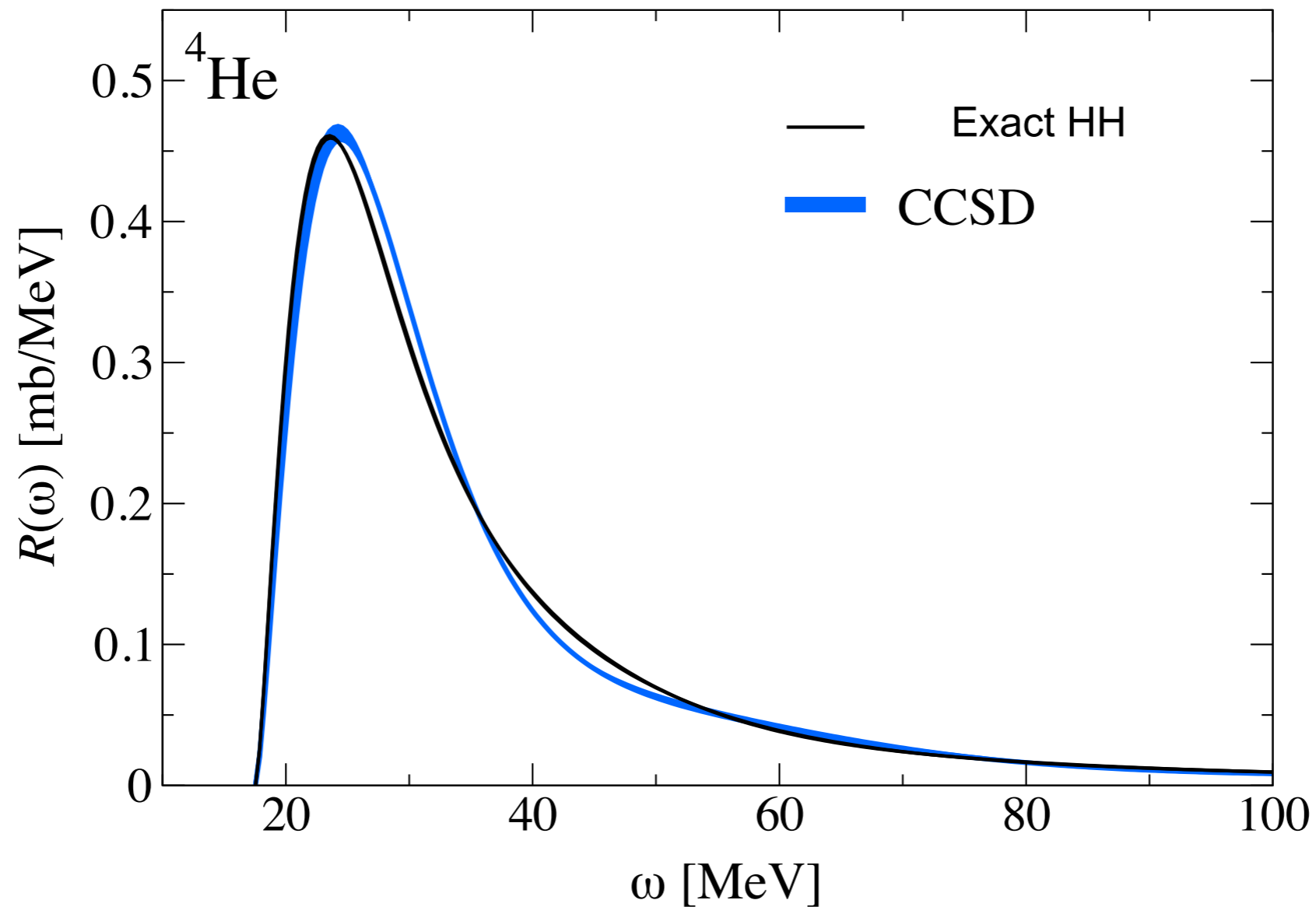
$\Gamma = 10$ [MeV]



S.B. *et al.*, Phys. Rev. Lett. **111**, 122502 (2013)

Dipole Response Functions with NN forces at N³LO

Validation ⁴He

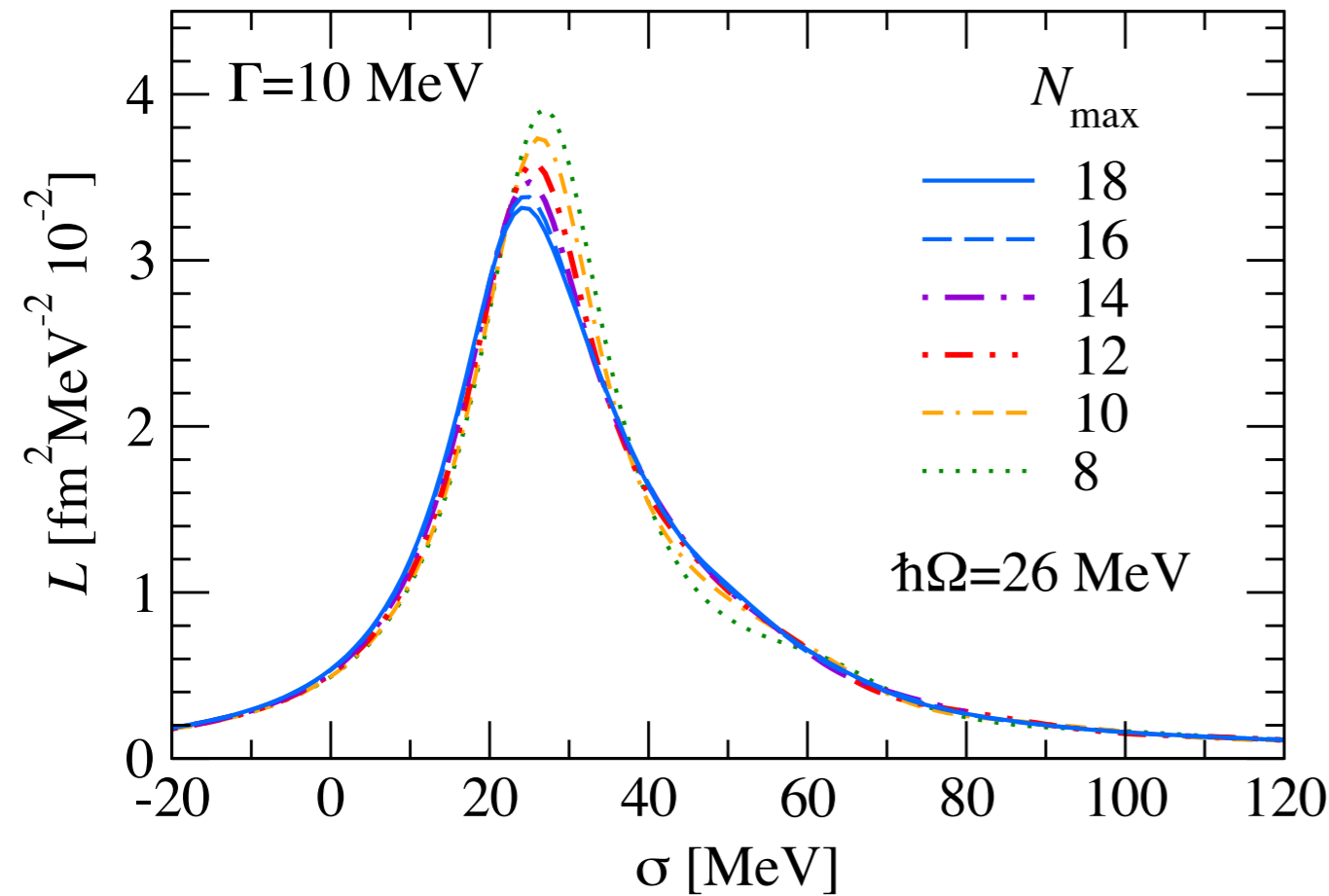


Pushing the mass limits

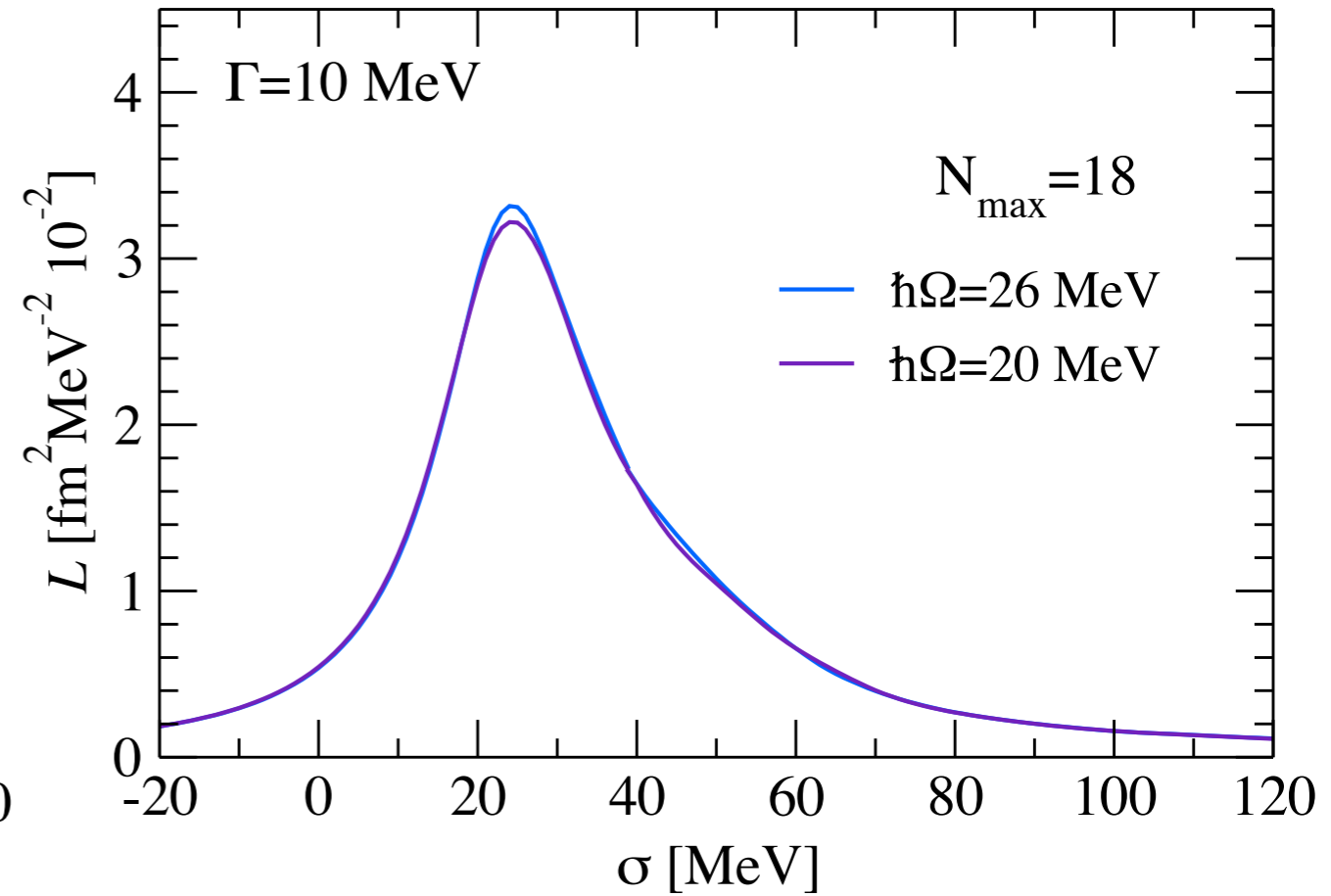
New theoretical method aimed at extending *ab-initio* calculations towards medium mass

Extension to Dipole Response Function in ^{16}O with NN forces derived from χEFT (N^3LO)

➔ Convergence in the model space expansion



Good convergence!



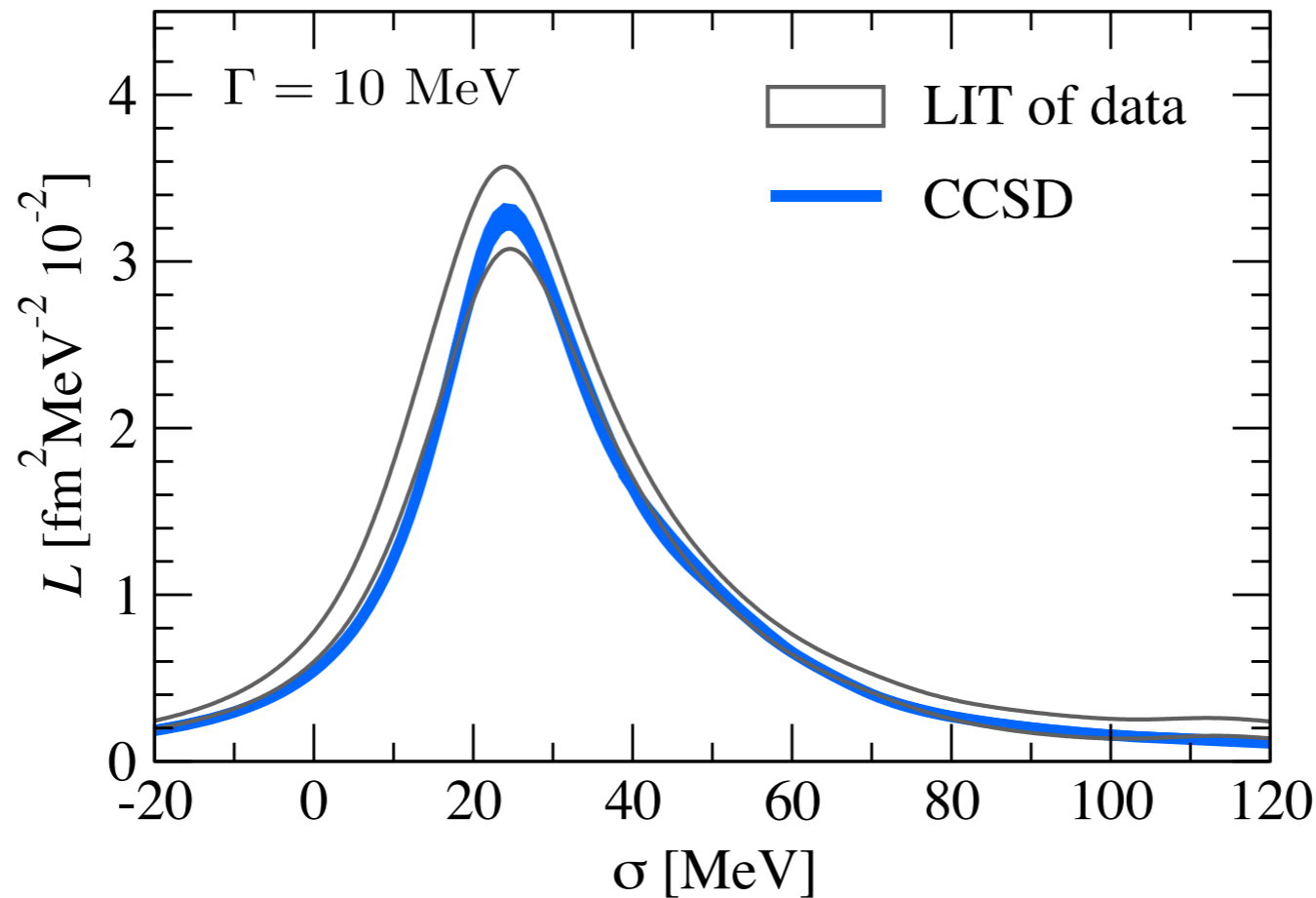
Small HO dependence: use it as error bar

Pushing the mass limits

New theoretical method aimed at extending *ab-initio* calculations towards medium mass

Extension to Dipole Response Function in ^{16}O with NN forces derived from χEFT (N^3LO)

➔ Comparison to the experiment

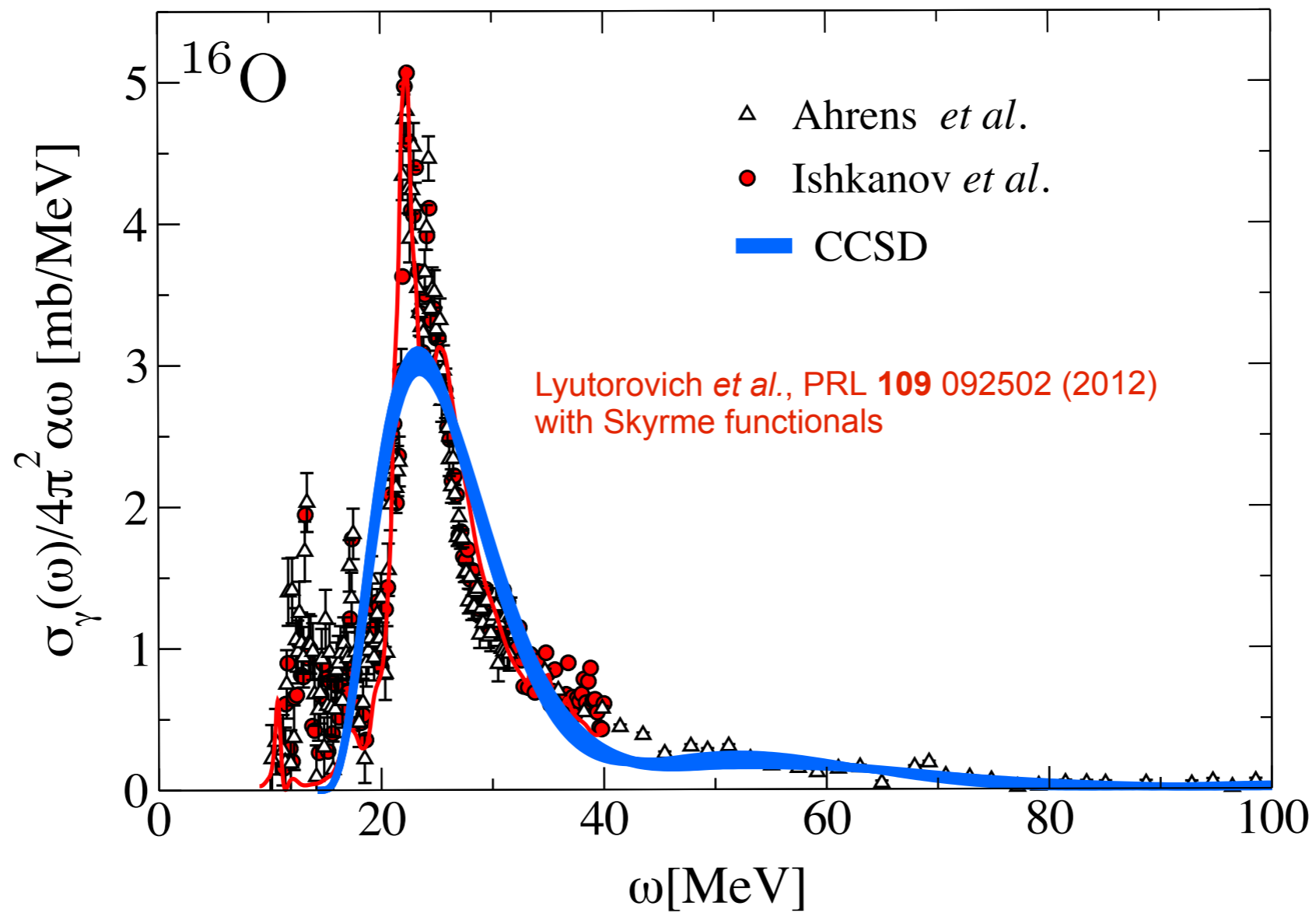


This proves that the total strength is correctly reproduced

Photonuclear reactions

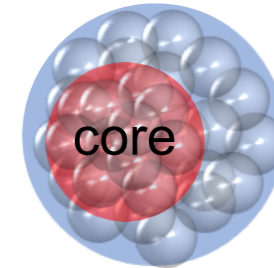
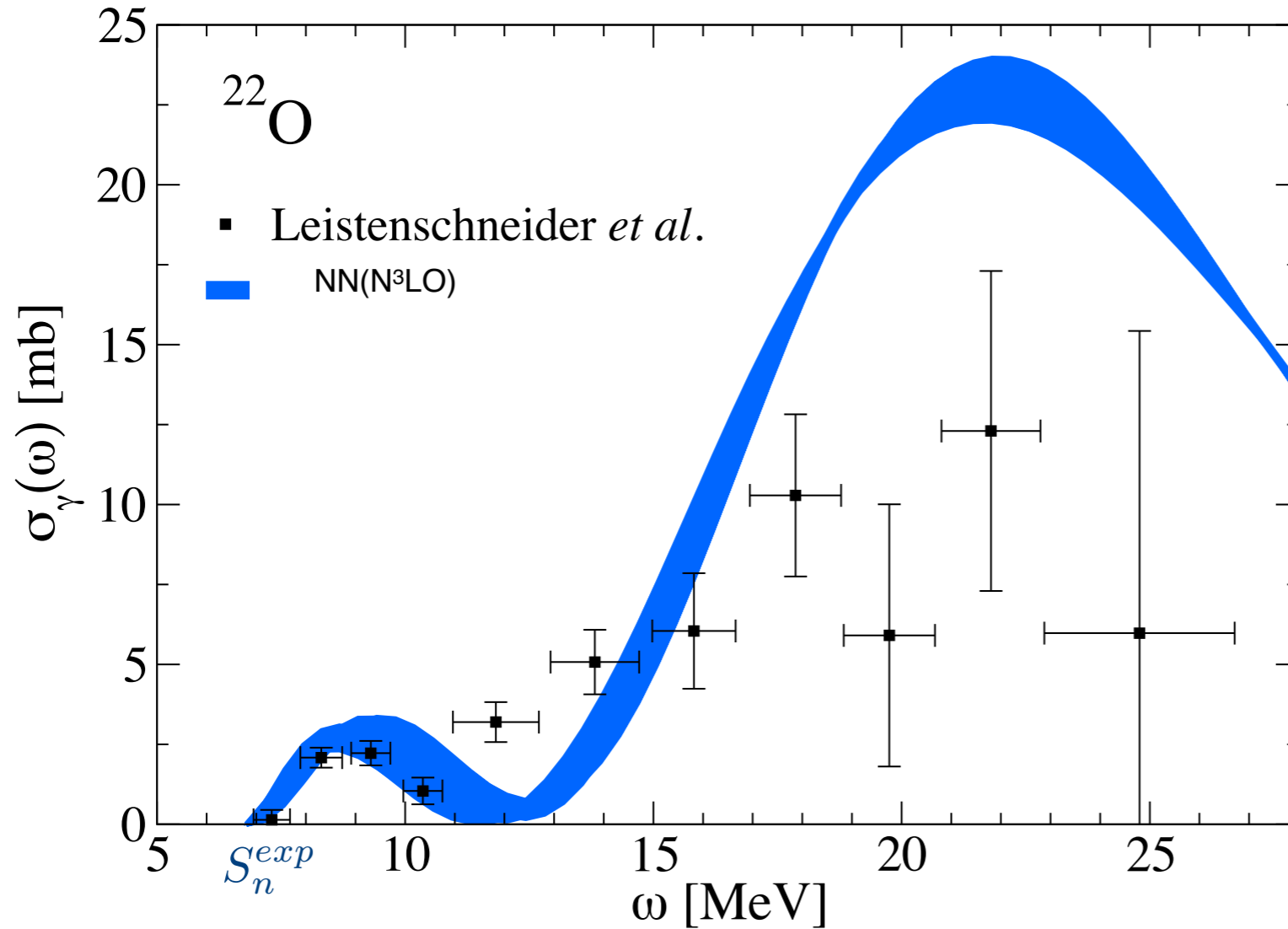
S.B. *et al.*, PRL **111**, 122502 (2013)

Photoabsorption cross section with NN forces at N³LO



Photoexcitation of neutron-rich nuclei

S.B. *et al.*, PRC **90**, 064619 (2014)

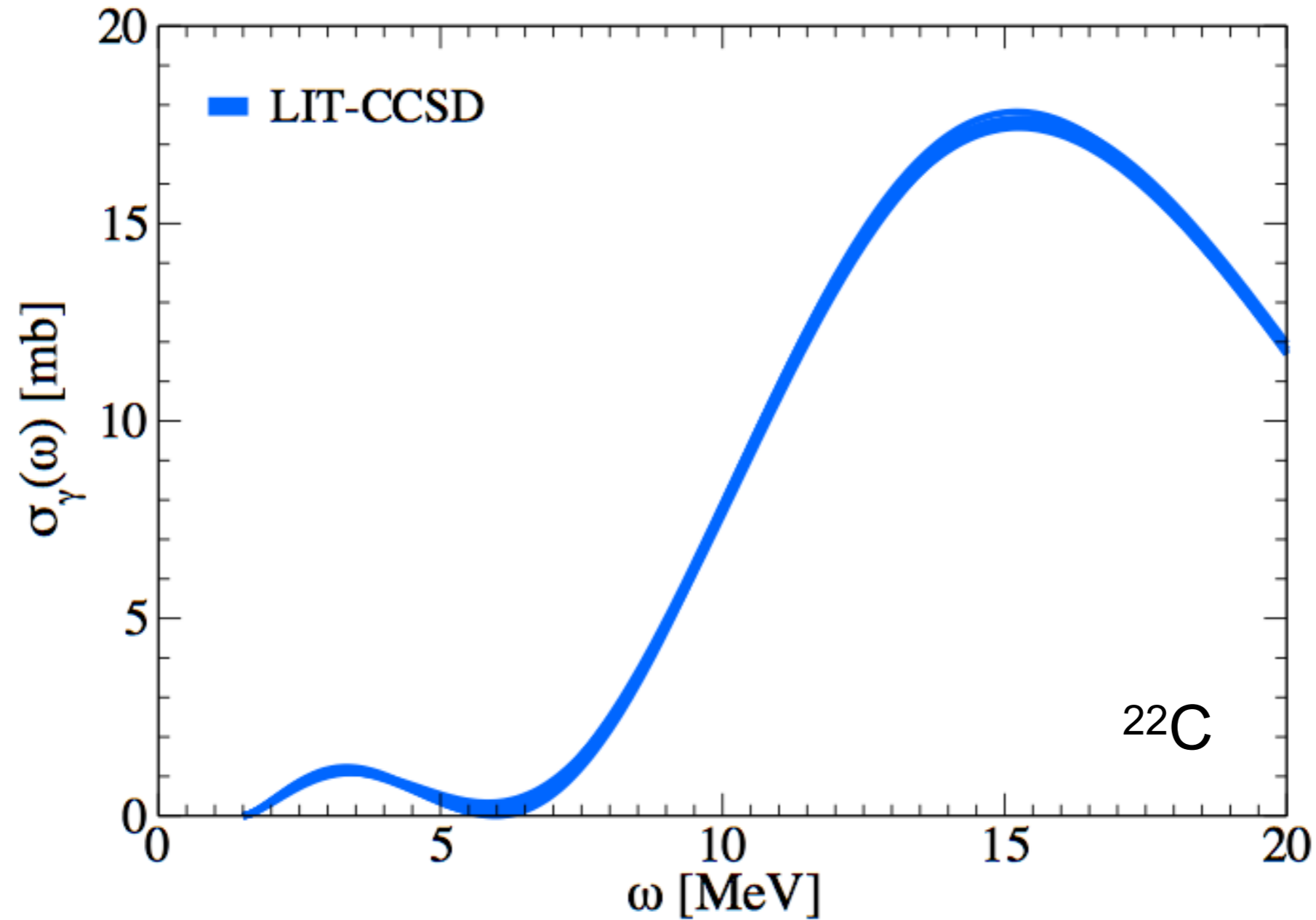


Pigmy Dipole Resonance (PDR)

Nicely described by a first principle calculation

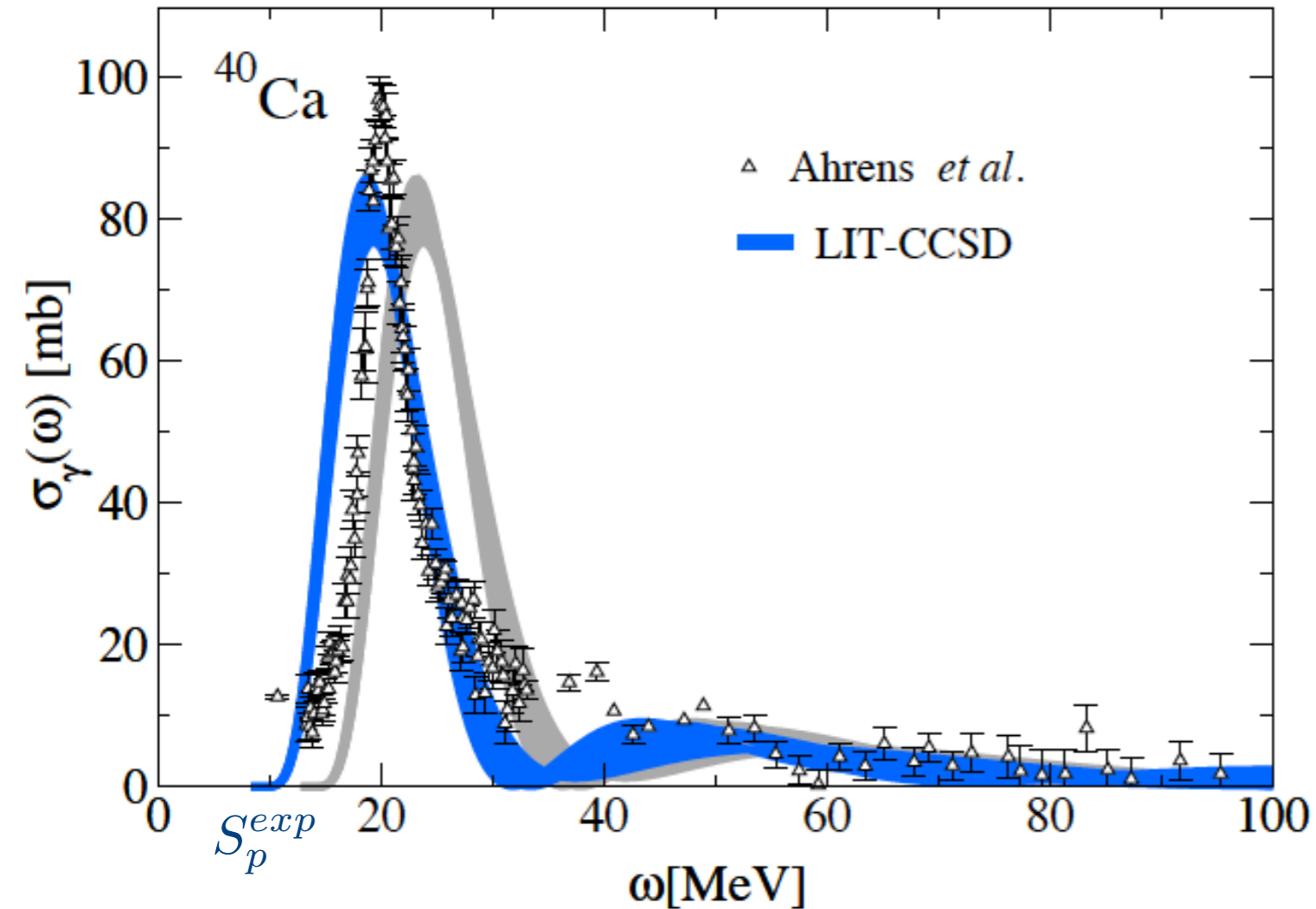
Theory provides a deeper understanding: microscopic interpretation of collective phenomena

M.Miorelli et al.



^{22}C being measured at RIKEN

Photonuclear reactions



Dipole resonance at too high energy related to too small electric Polarizability and too small radii

$$\alpha_D \rightarrow \int_0^{E_x} d\omega \frac{\sigma_\gamma(\omega)}{\omega^2}$$

$$\alpha_D(\text{Th}) = 1.47 \text{ fm}^3$$

$$\alpha_D(\text{Exp}) = 2.23(3) \text{ fm}^3$$

$$R_{\text{ch}}(\text{Th}) = 3.05 \text{ fm}$$

$$R_{\text{ch}}(\text{Exp}) = 3.48 \text{ fm}$$

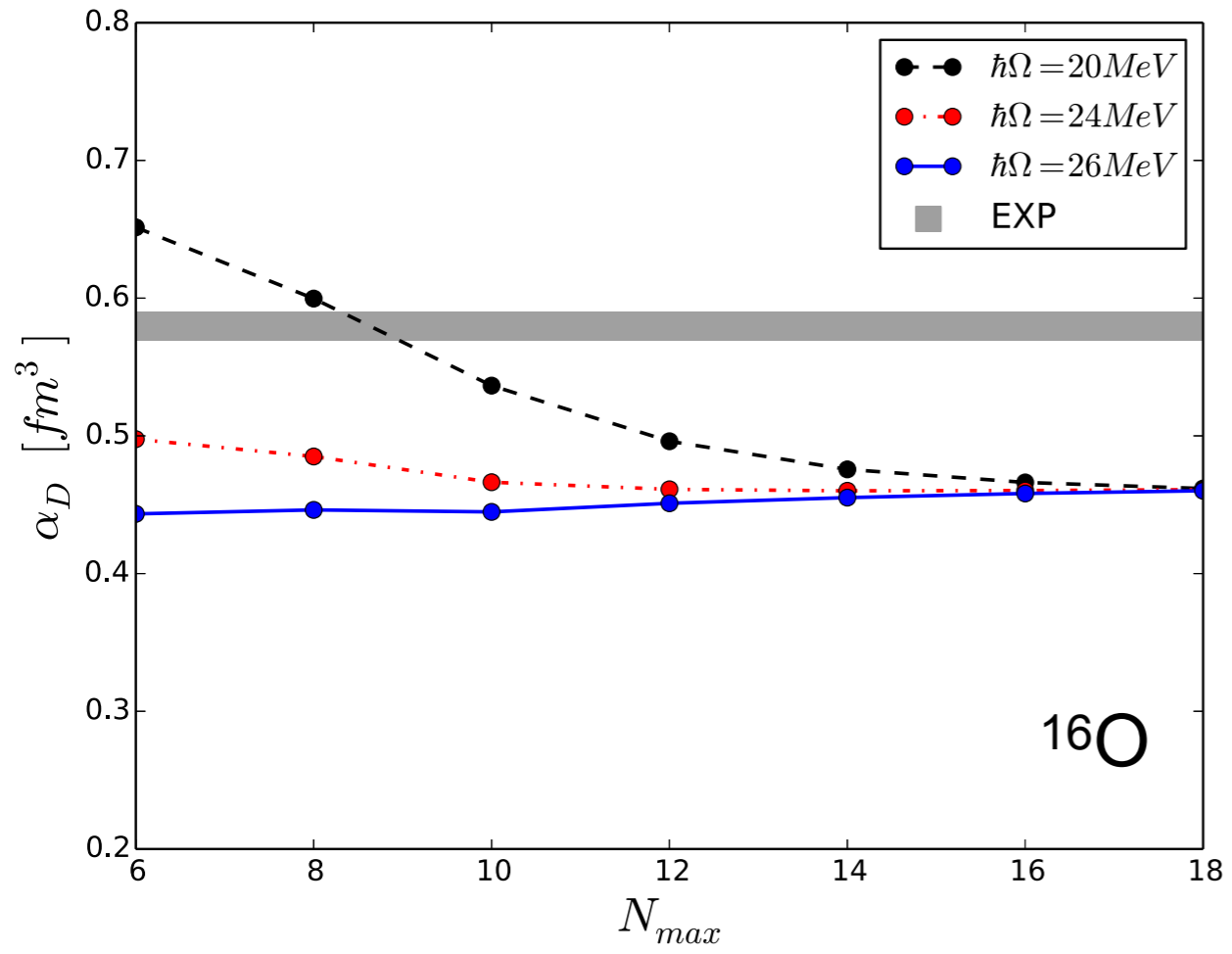
$$E(\text{Th}) = 362 \text{ MeV}$$

$$E(\text{Exp}) = 342 \text{ MeV}$$

Radii and binding energies are off the experimental values, as well as dipole polarizability

Electric dipole polarizability

Medium-mass nuclei with NN(N³LO)

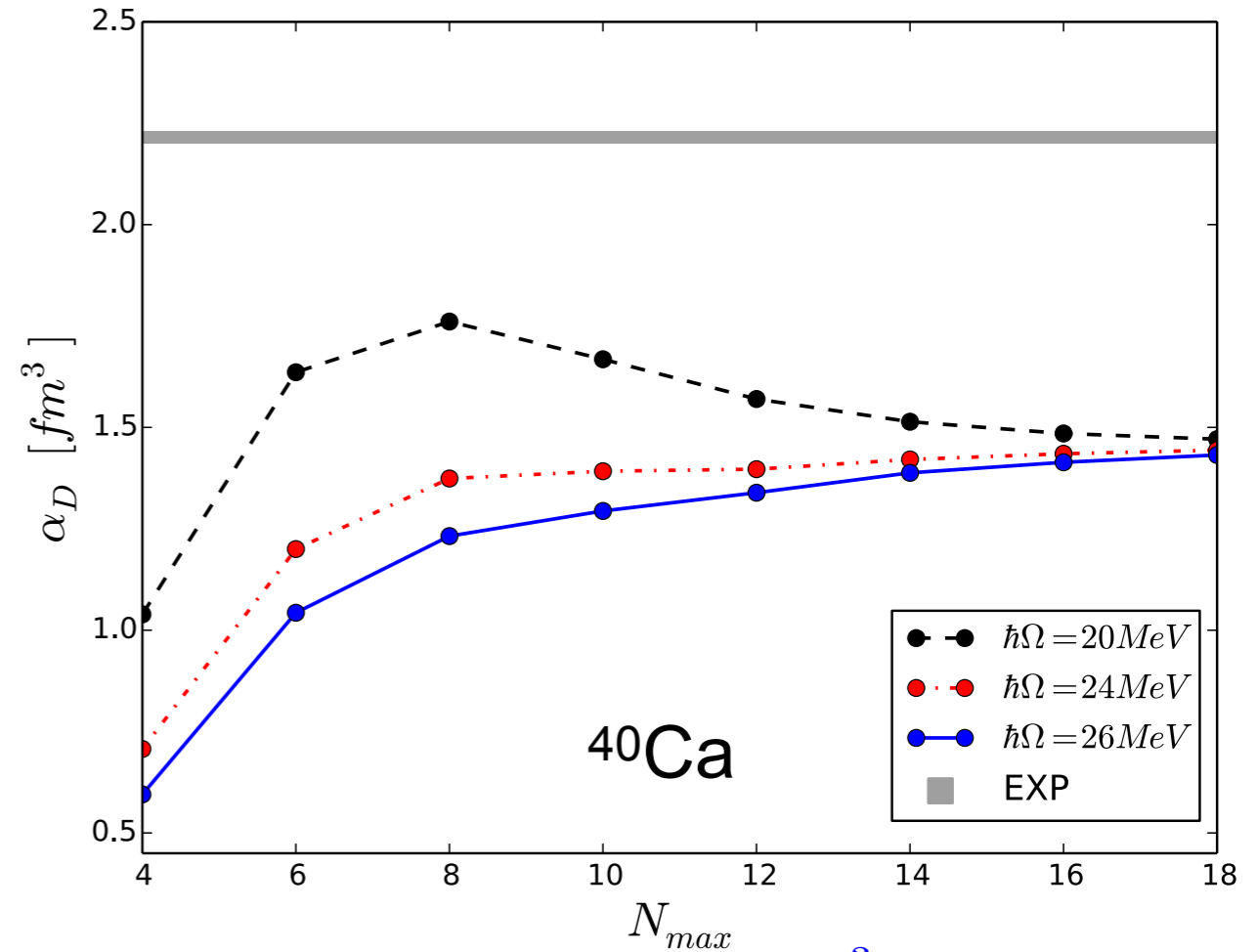


$$\alpha_D = 0.46 \text{ fm}^3$$

$$\alpha_D^{exp} = 0.585(9) \text{ fm}^3$$

$$R_{ch} = 2.3 \text{ fm}$$

$$R_{ch}^{exp} = 2.6991(52) \text{ fm}$$



$$\alpha_D = 1.47 \text{ fm}^3$$

$$\alpha_D^{exp} = 2.23(3) \text{ fm}^3$$

$$R_{ch} = 3.05 \text{ fm}$$

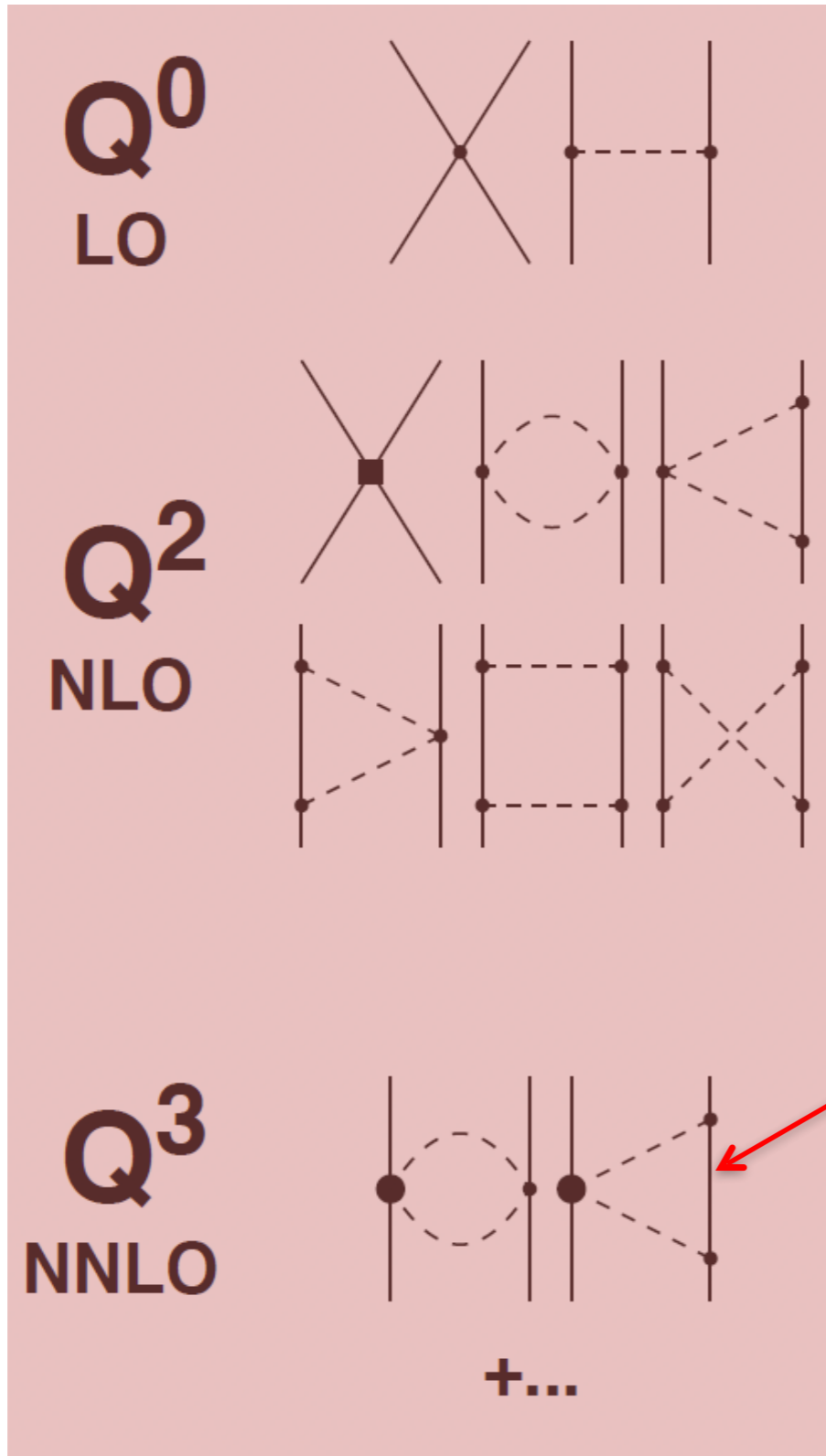
$$R_{ch}^{exp} = 3.4776(19) \text{ fm}$$

The present Hamiltonian underestimates both radii and electric dipole polarizabilities

What shall we do to solve it?

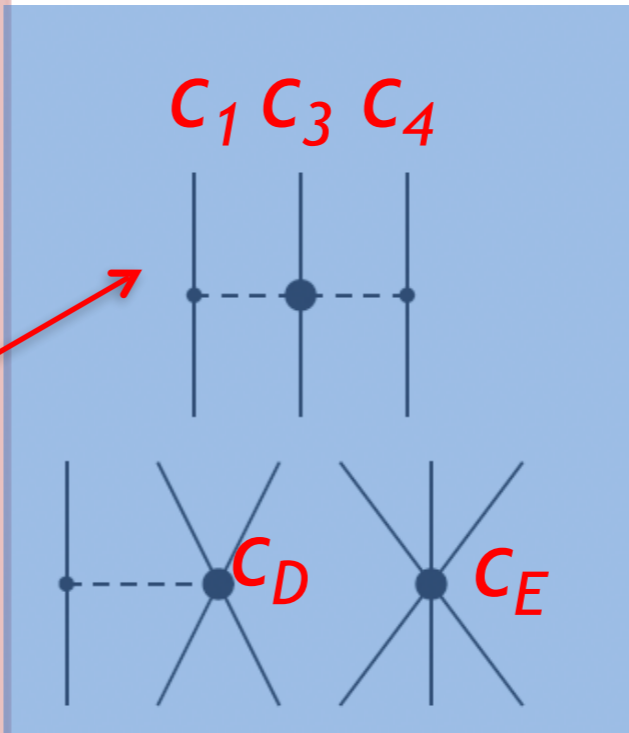
2N Force

3N Force



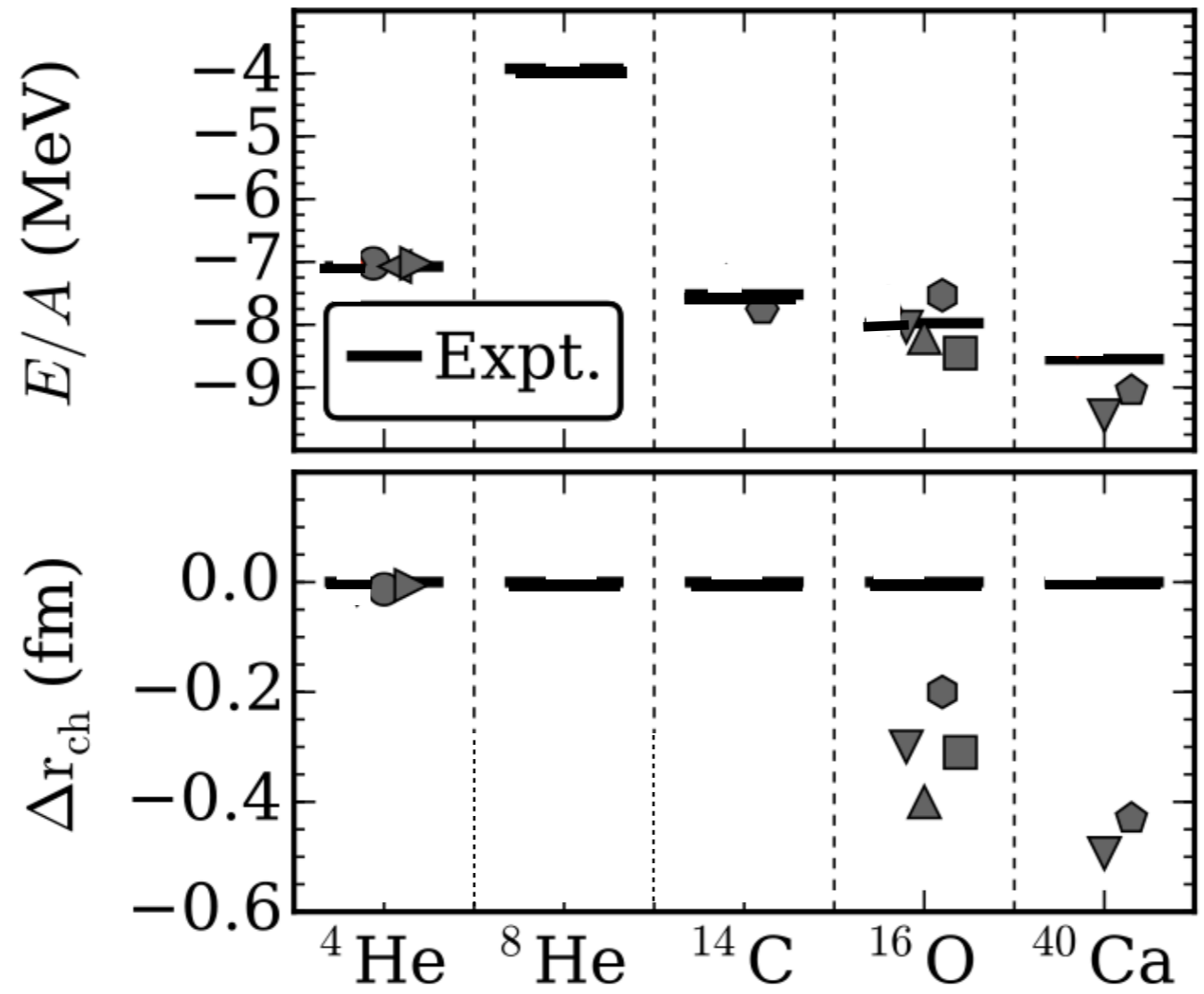
Traditional approach:

- Fit interactions nucleon by nucleon
- Fit to NN scattering data up to $\sim 350\text{MeV}$
- cE and cD fit to $A=3,4$



Add three-nucleon forces

That does not always work

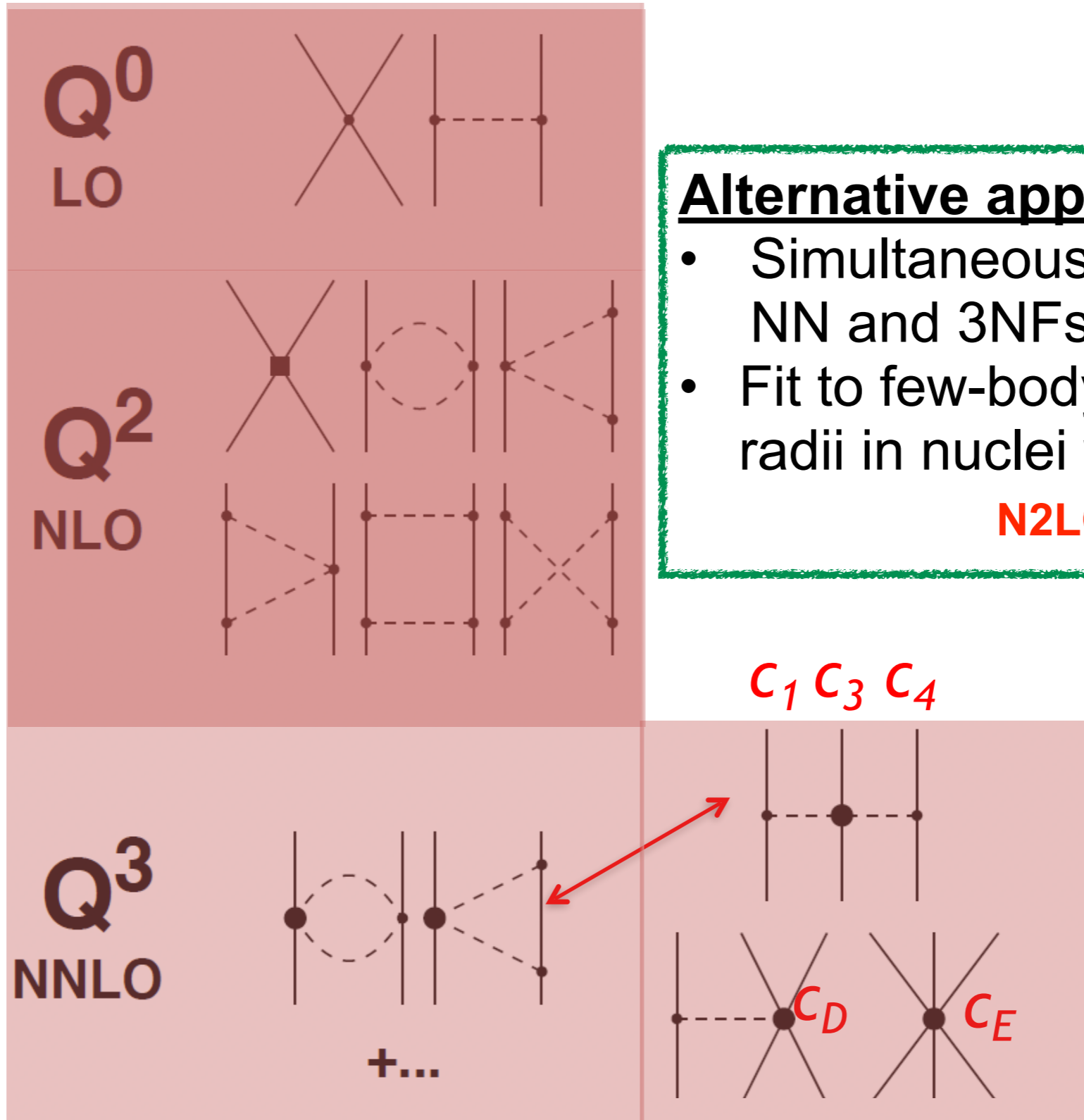


Alternative strategy

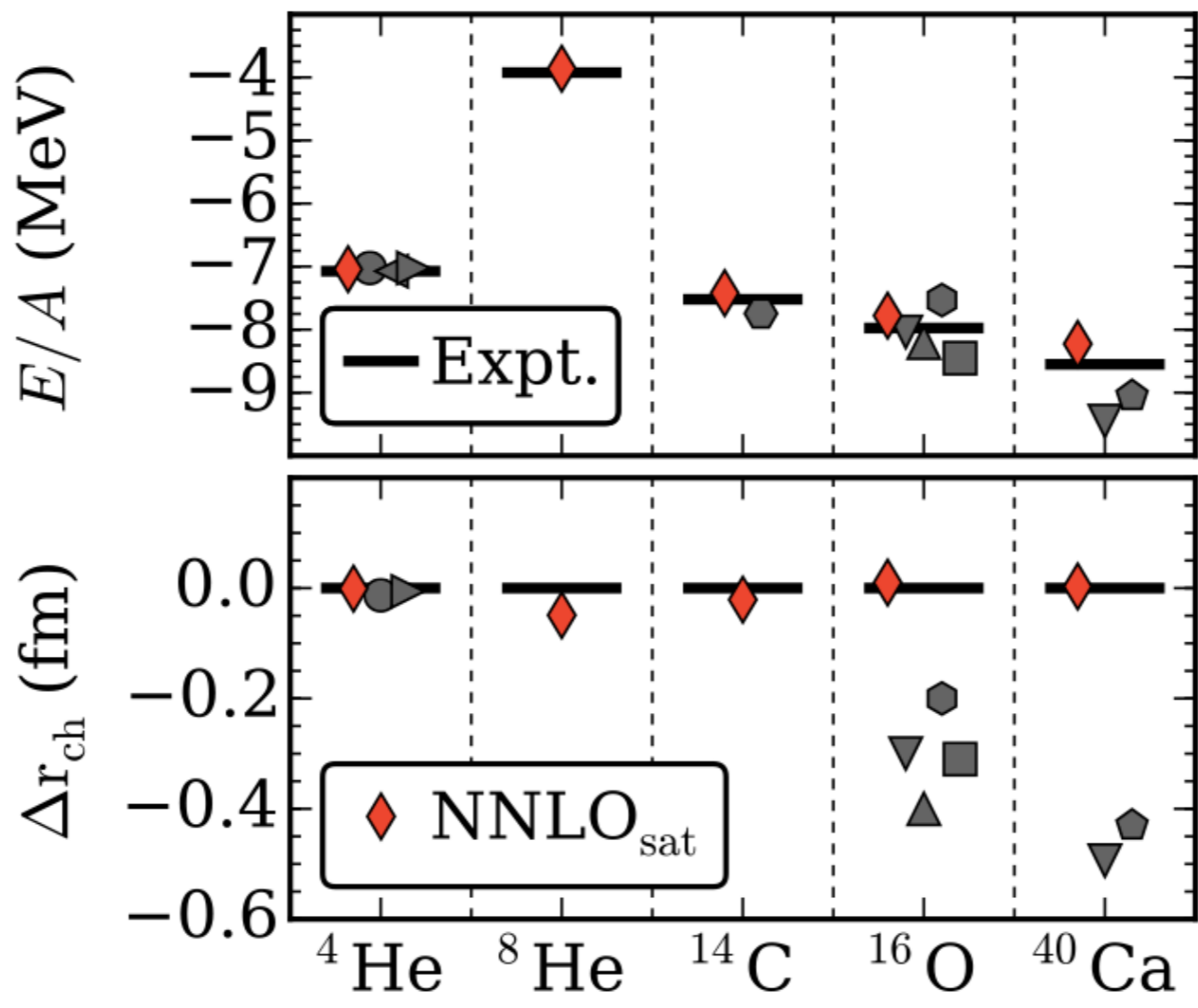
A. Ekström *et al*, Phys. Rev. C 91, 051301(R) (2015)

2N Force

3N Force



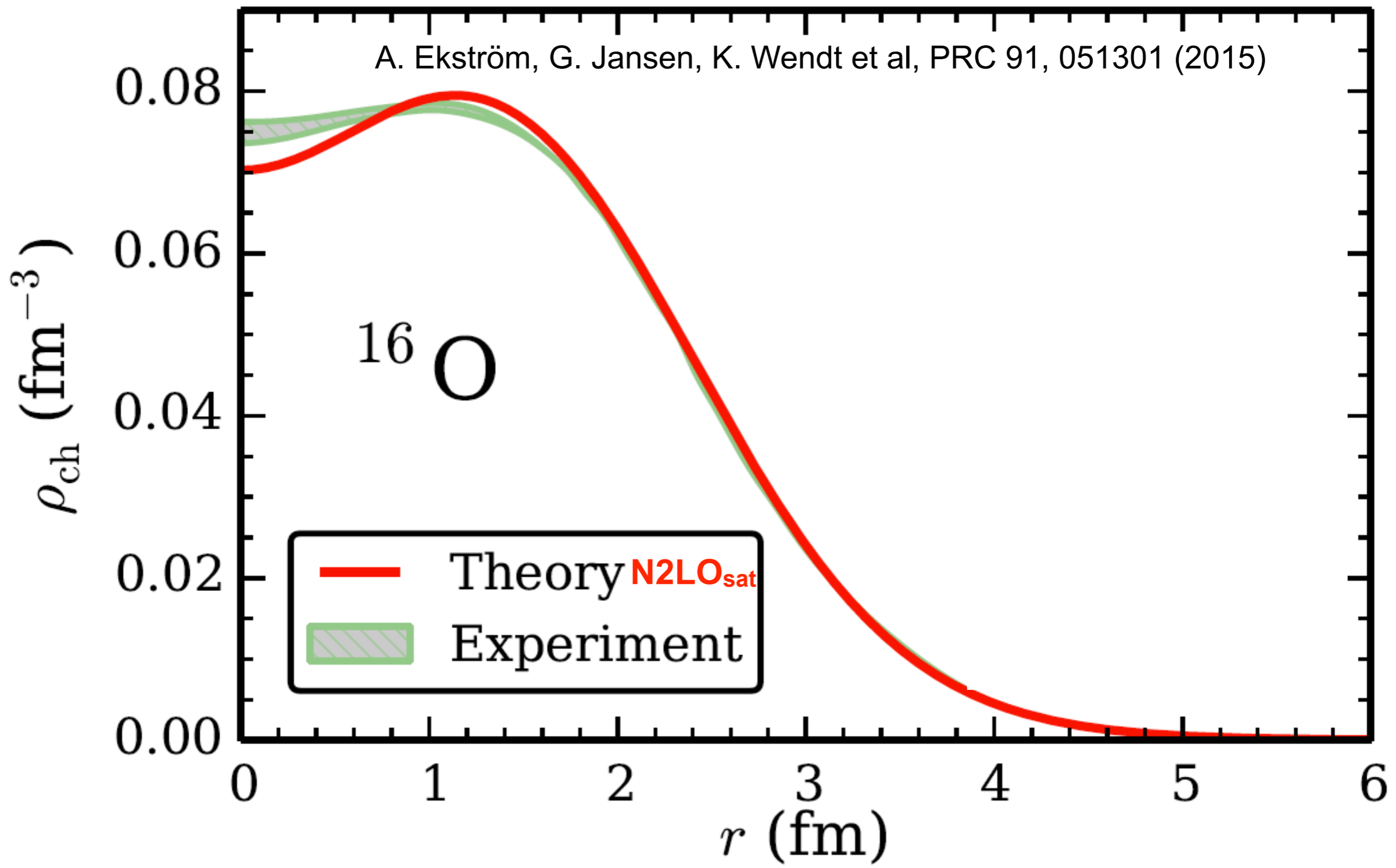
Fixes the radii



Include radii in the fit of LEC for the three-body force
 ^{40}Ca is a pure prediction

Charge density of ^{16}O

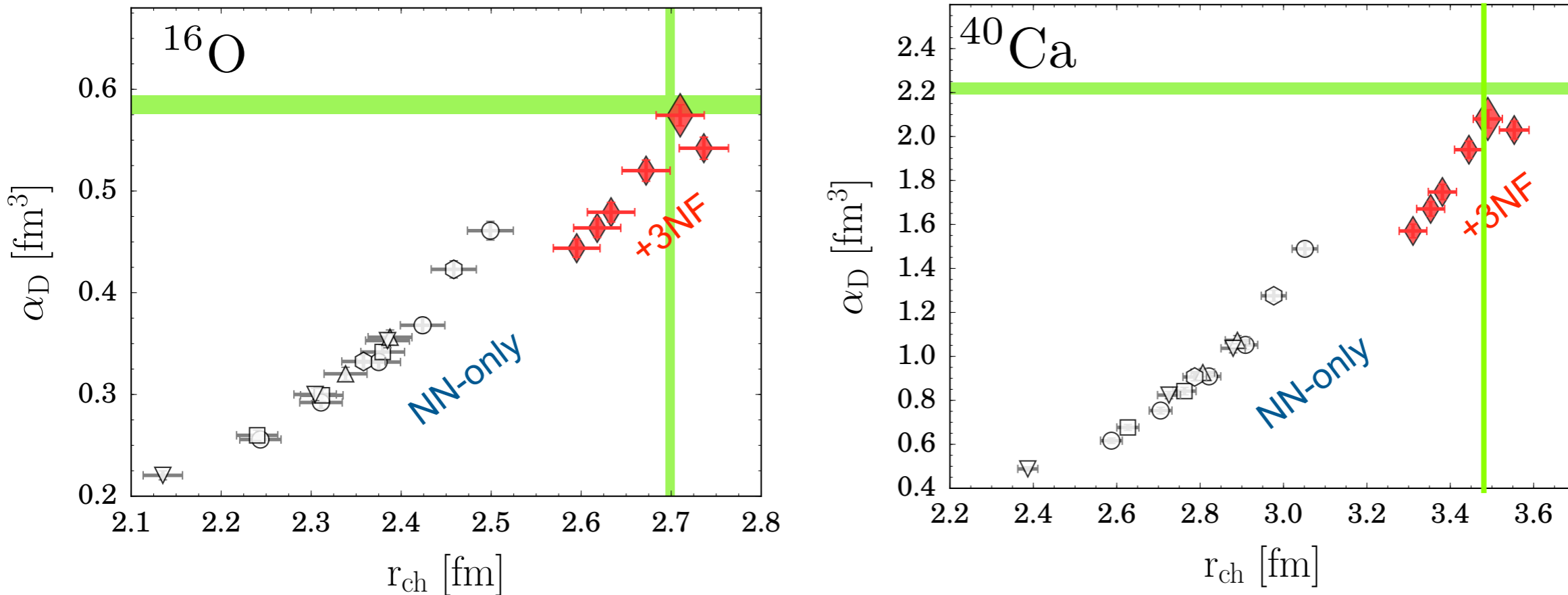
A. Ekström, G. Jansen, K. Wendt et al, PRC 91, 051301 (2015)



Electric Dipole Polarizability

Medium-mass nuclei with NN + 3NF interactions

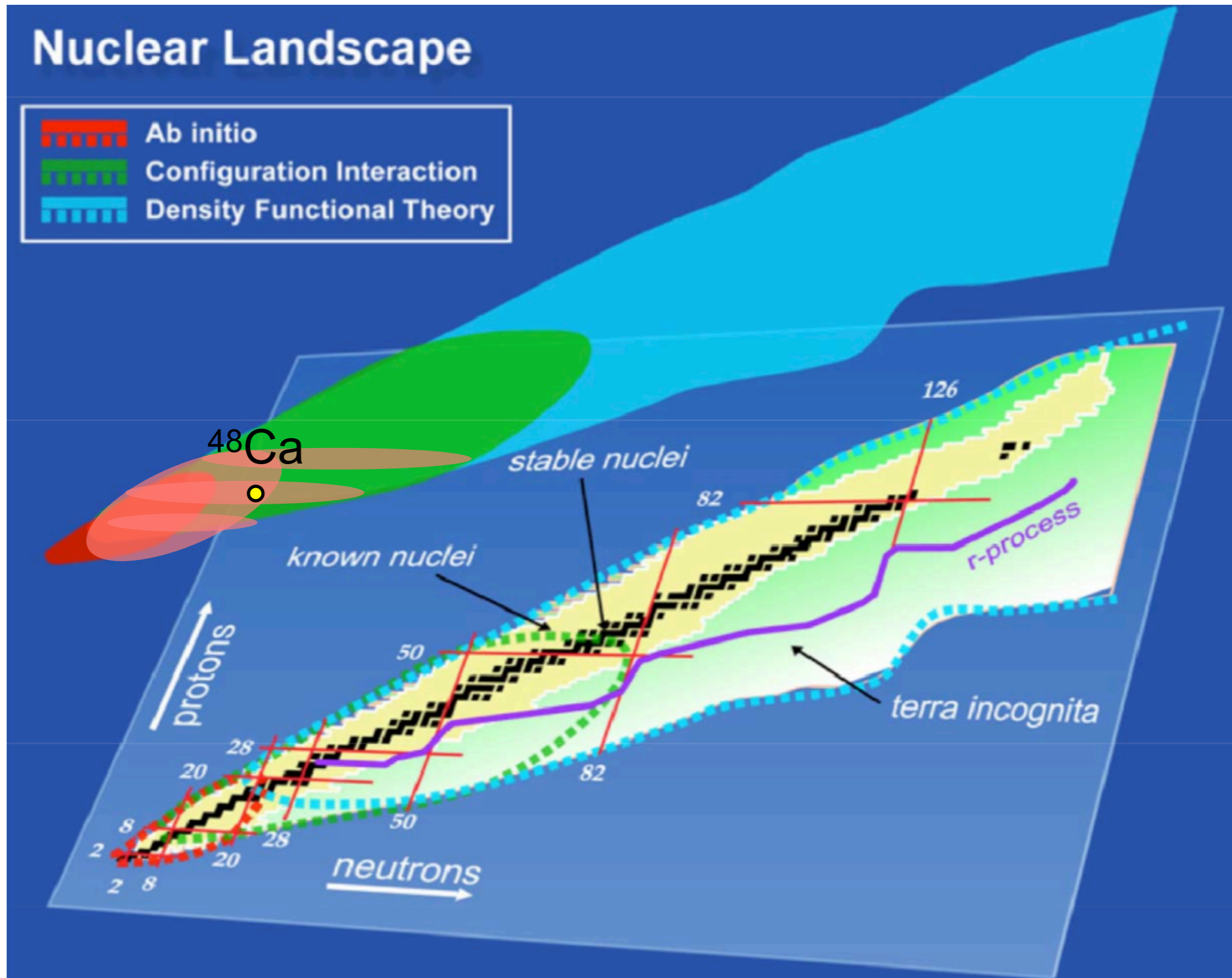
M. Miorelli *et al.*, PRC **94** 034317 (2016)



3NF A. Ekström *et al.*, Phys. Rev. C91, 051301 (2015)
 K. Hebeler *et al.*, Phys. Rev. C83, 031301 (2011)

Much better agreement with experimental data
 Variation of Hamiltonian can be used to assess the theoretical error bar

What about heavier systems?



<http://unedf.org>

While neutron-rich, for all practical purposes it can be considered a stable nucleus

★ (p,p') scattering to extract the **electric dipole polarizability** at RCNP, Japan

α_D is related to the symmetry energy in the EOS of nuclear matter

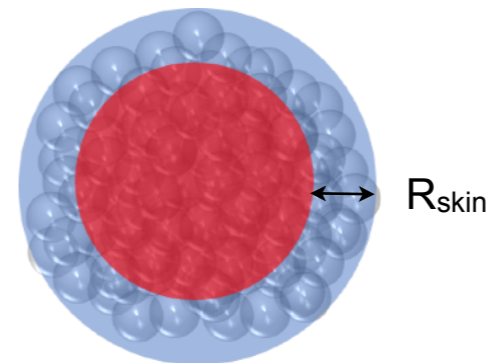
★ **Parity violation electron scattering** Calcium Radius Experiment (CREX) at JLab and the Mainz Radius Experiment (MREX) at MAMI/MESA to measure R_{skin}

$$A_{pv} = \frac{d\sigma/d\Omega_R - d\sigma/d\Omega_L}{d\sigma/d\Omega_R + d\sigma/d\Omega_L} \approx -\frac{G_F q^2}{4\pi\alpha\sqrt{2}} \frac{Q_W F_W(q^2)}{Z F_{ch}(q^2)}$$

The weak force probes the neutron distribution

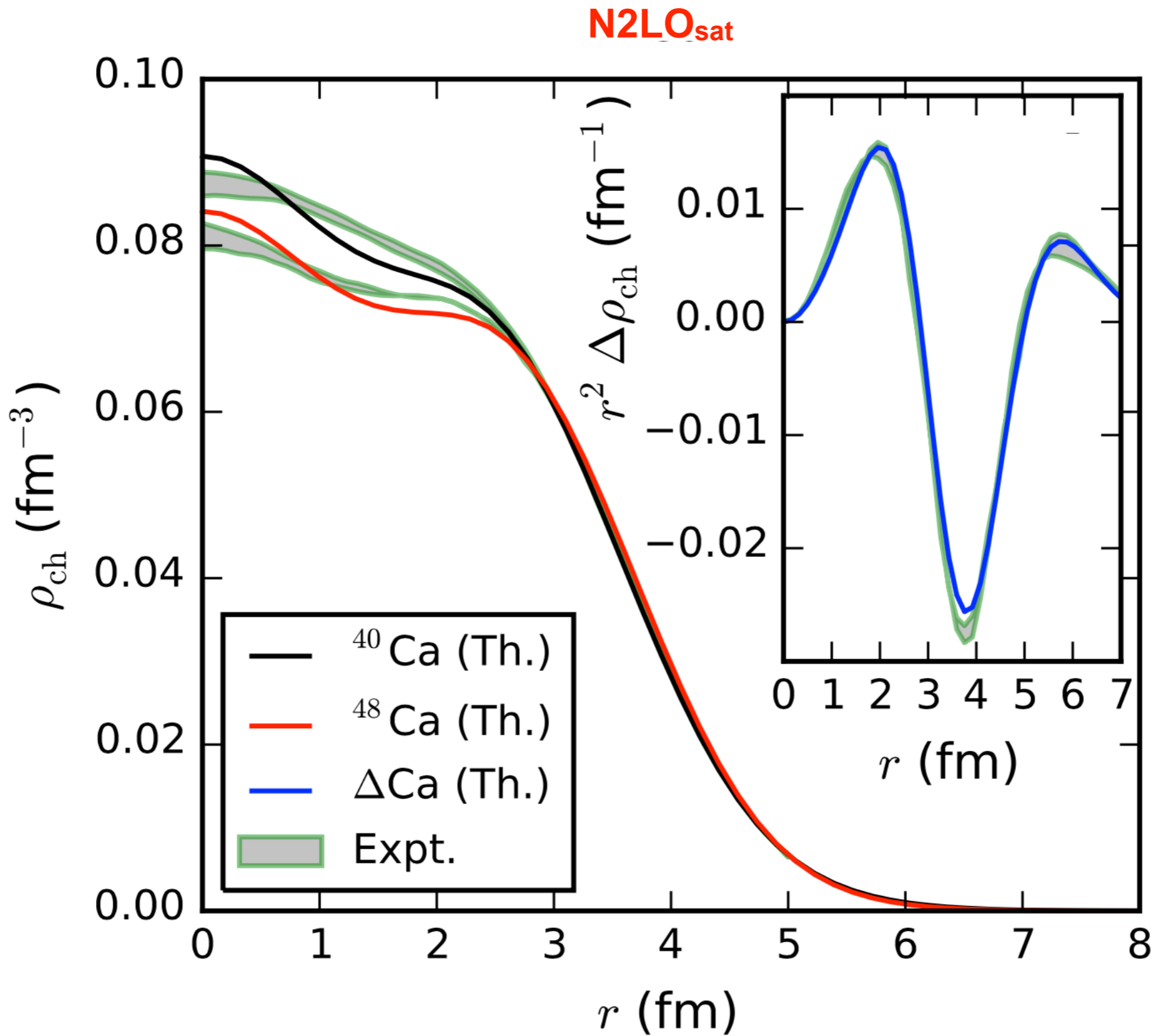
$$Q_W^n \approx -1$$

$$Q_W^p = 1 - 4 \sin^2 \theta_W \approx 0$$



Can we give a first principle predictions for these future experiments?

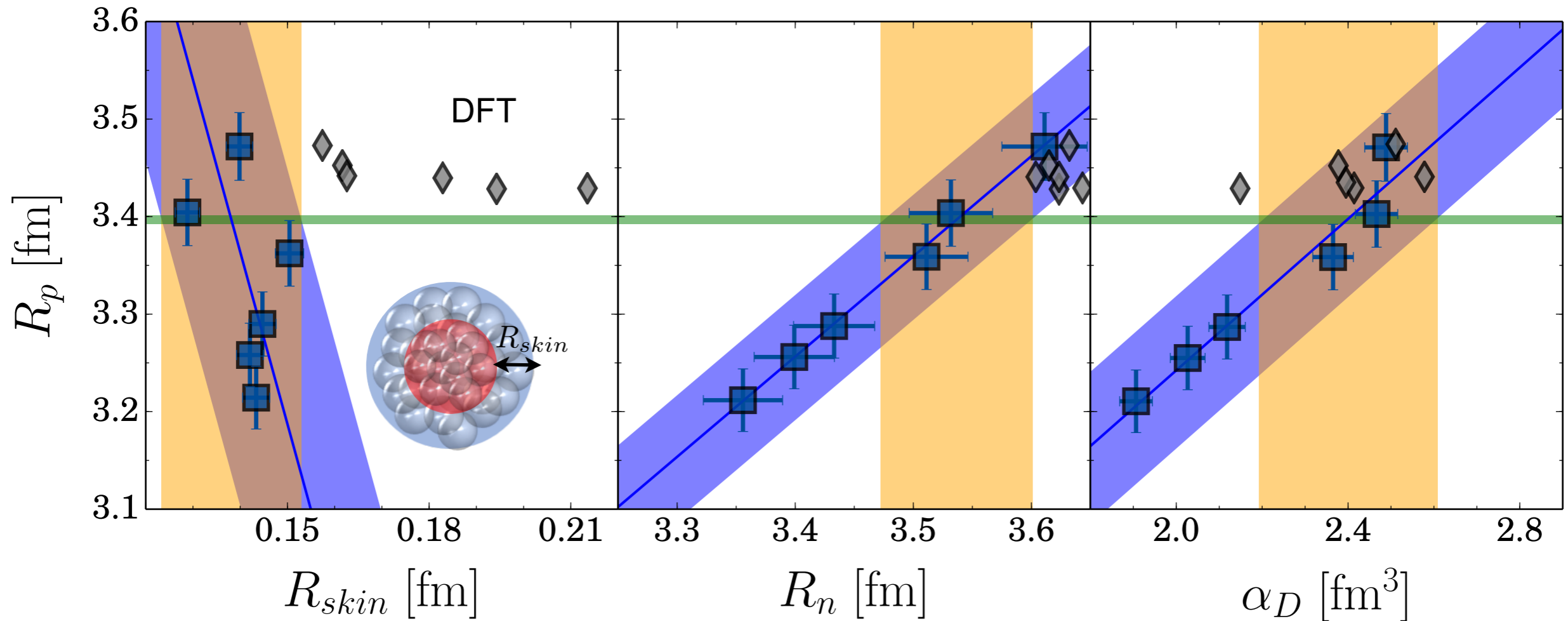
Charge density of ^{48}Ca



^{48}Ca from first principles

International collaboration (USA/Canada/Europe/Israel) using coupled-cluster theory

Hagen *et al.*, *Nature Physics* **12**, 186 (2016)



■ Ab initio with **three nucleon forces** from chiral EFT

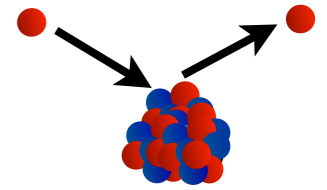
Strong correlations with R_p allow to put narrow constraints to R_{skin} and α_D

Ab-initio predictions: $0.12 \leq R_{skin} \leq 0.15 \text{ fm}$
 $2.19 \leq \alpha_D \leq 2.60 \text{ fm}^3$

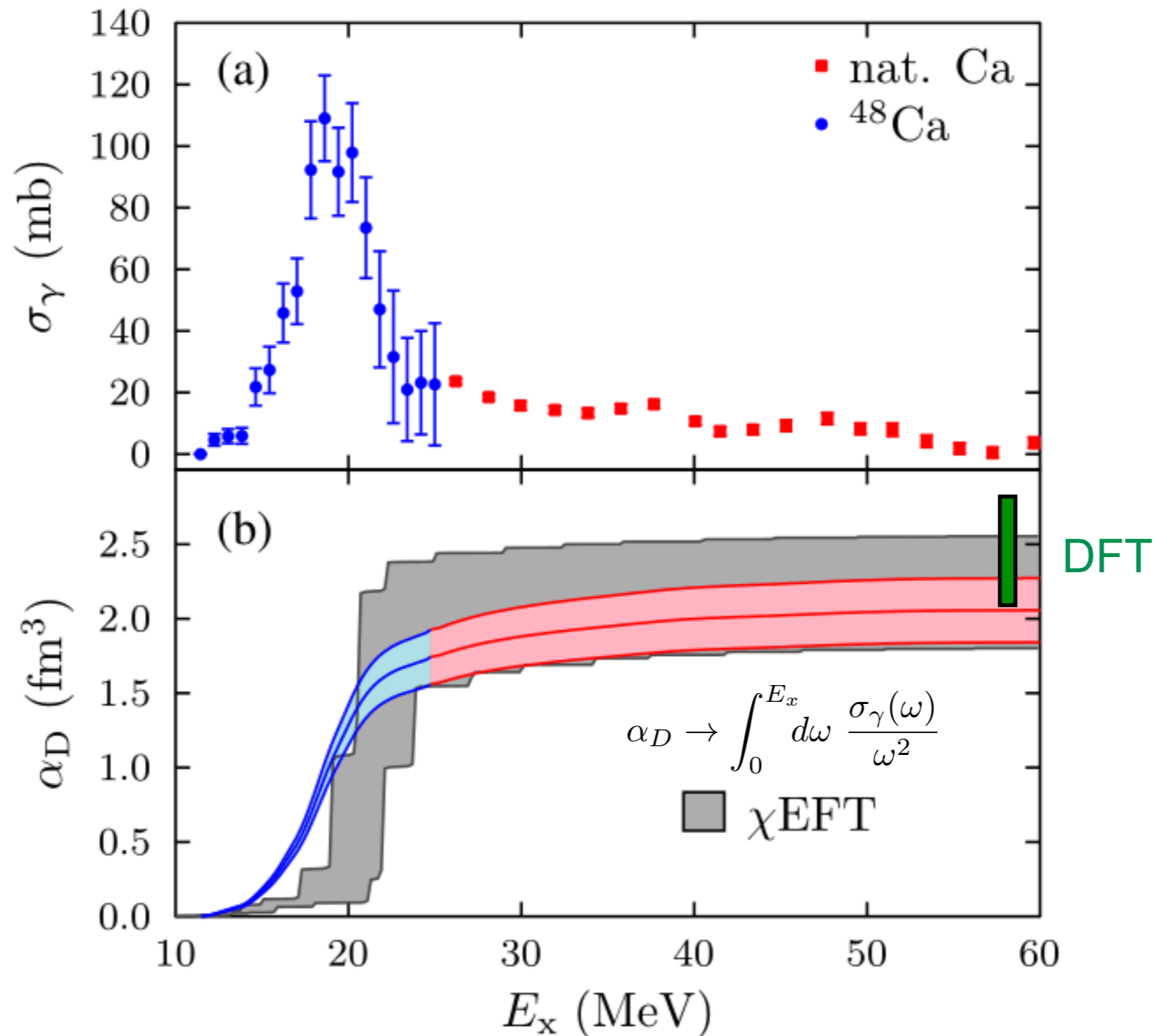
R_{skin} will be measured with **Parity violating electron scattering** CREX MREX

^{48}Ca electric dipole polarizability

New measurements from the Osaka-Darmstadt collaboration using inelastic proton scattering



J.Birkhan, *et al.*, Phys. Rev. Lett. **118**, 252501 (2017)



S.B. *et al.*, PRL **111**, 122502 (2013)

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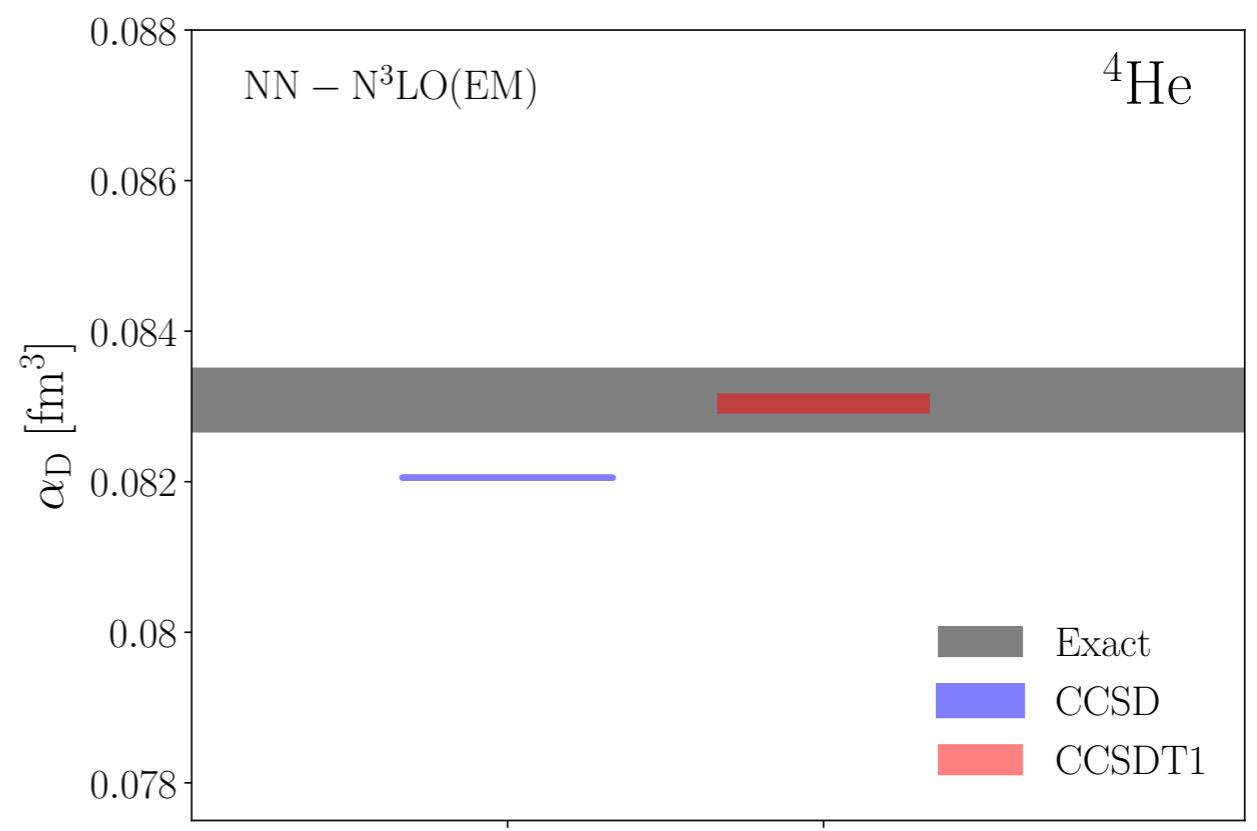
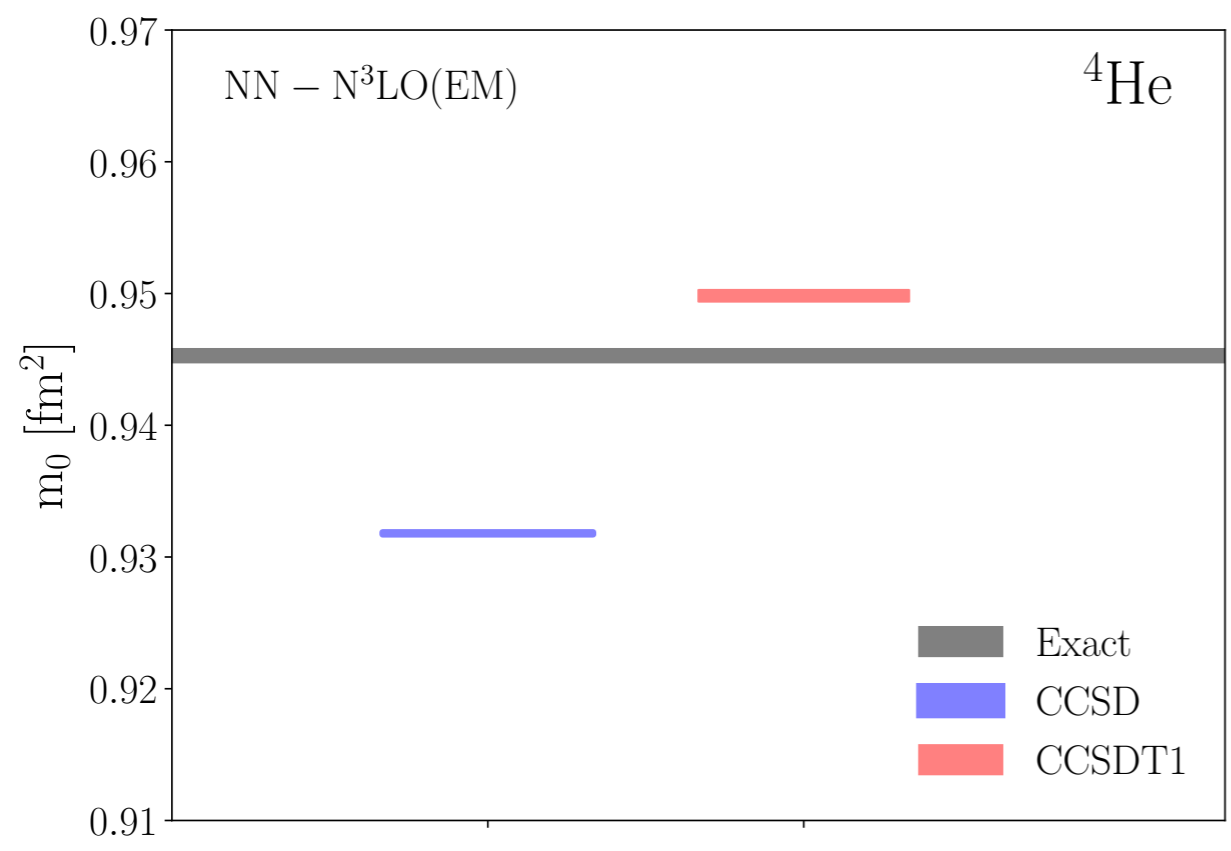
$$(\bar{H} - z^*)|\tilde{\Psi}_R(z^*)\rangle = \bar{\Theta}|\Phi_0\rangle$$

The first implementation we provided is based on the CCSD approximation scheme

Look at triples corrections

How to improve our tools

M. Miorelli *et al.*, in preparation (2017)



CCSD scheme $e^T = e^{T_1+T_2}$

$$R = R_0 + R_1 + R_2$$

CCSDT1 scheme (linearized triples)

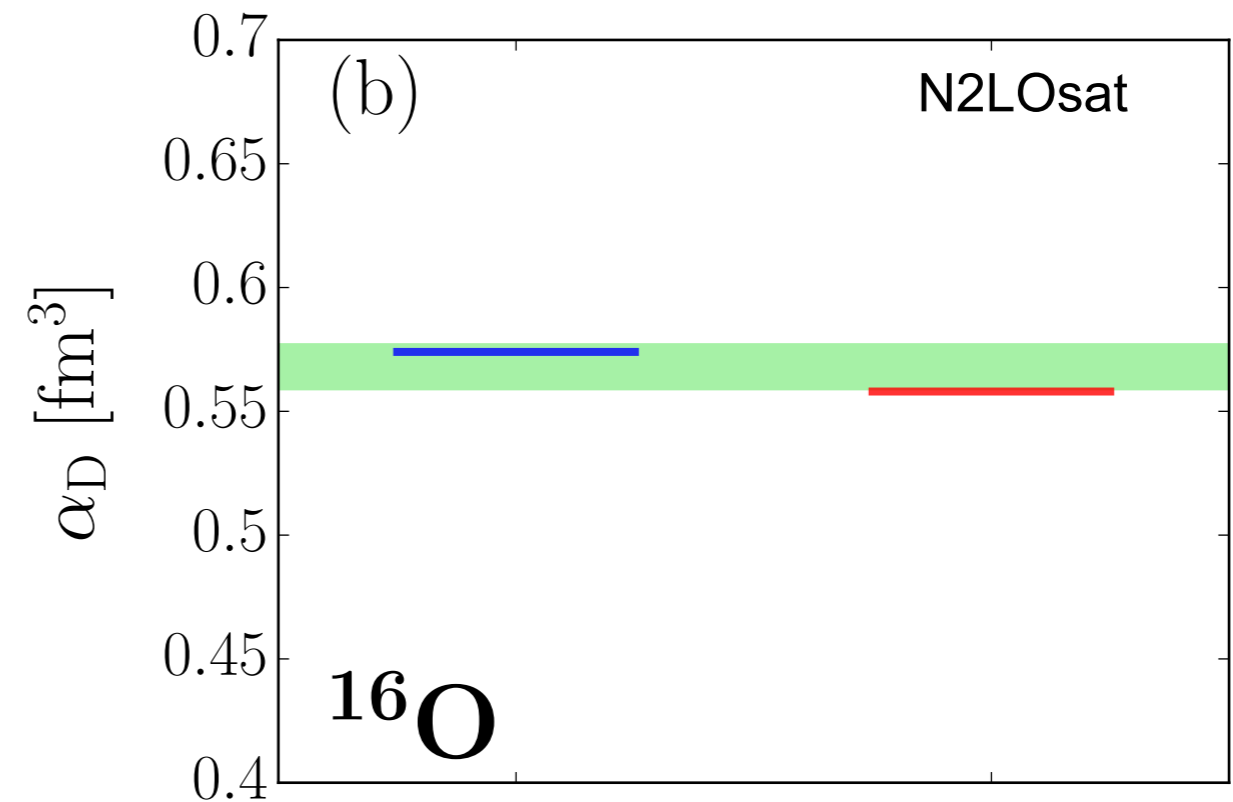
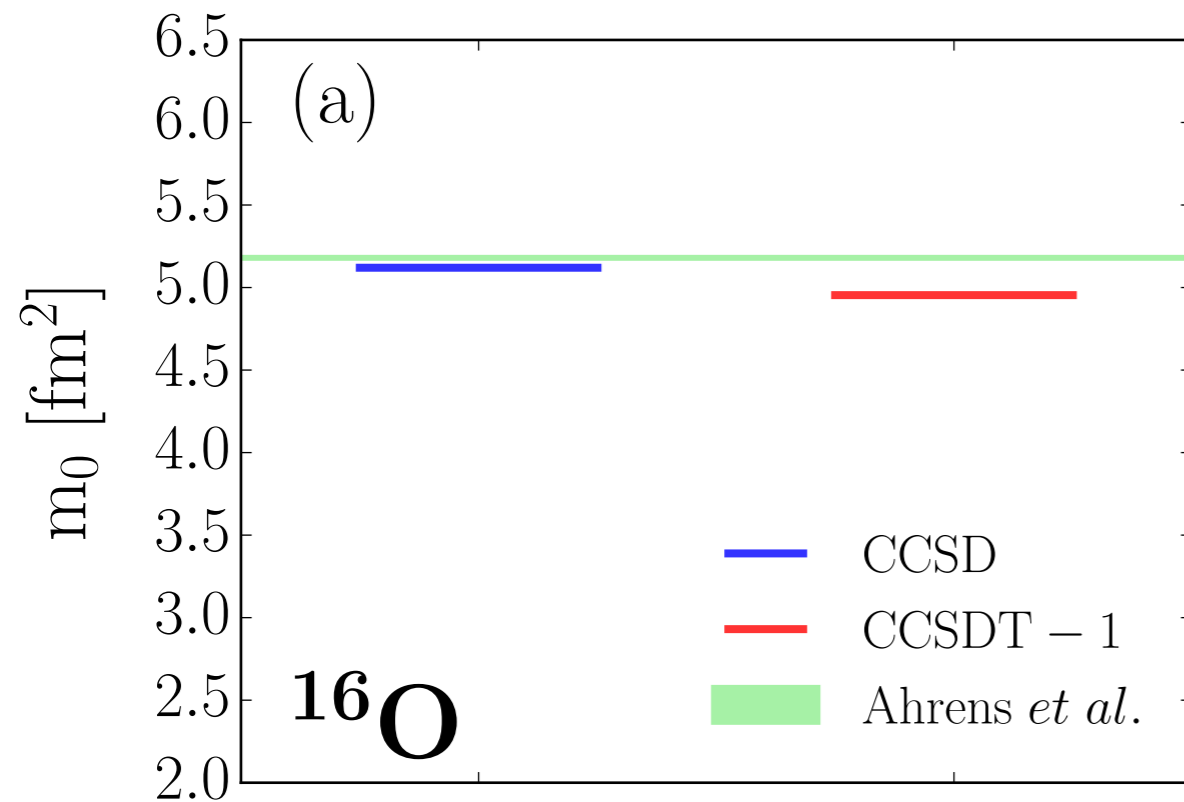
$$e^T = e^{T_1+T_2} + T_3$$

$$R = R_0 + R_1 + R_2 + R_3$$

Exact \Rightarrow hyperspherical harmonics, all correlations included (up to quadruples)

How to improve our calculations

M. Miorelli *et al.*, in preparation (2017)



CCSD scheme $e^T = e^{T_1+T_2}$

$$R = R_0 + R_1 + R_2$$

CCSDT1 scheme
(linearized triples) $e^T = e^{T_1+T_2} + T_3$

$$R = R_0 + R_1 + R_2 + R_3$$

More on experimental techniques from *Christopher Lehr and Philipp Ries*

More on theoretical techniques from *Johannes Simonis*