Non-observability and reaction-structure interplay

Topical lecture week: Nuclear Structure from Spectroscopy and Direct Reactions

Kai Hebeler Darmstadt, Feb 24, 2016









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Reaction processes within ab initio frameworks



Calculations generally involve:

- nuclear structure part for description of initial/final state wave functions
- reaction part describes interaction with external probes

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- nuclear structure part for description of initial/final state wave functions
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For calculations of observables a factorization of structure and reaction parts has to be assumed:

$$\sigma \sim \langle \psi_F(\lambda) | O(\lambda) | \psi_A(\lambda) \rangle$$

whereas λ is a chosen resolution scale

"Traditional" NN interactions



- constructed to fit scattering data (long-wavelength information)
- long-range part dominated by one pion exchange interaction
- short range part strongly model dependent!
- traditional NN interactions contain strongly repulsive core at small distance
 - strong coupling between low and high-momenta
 - many-body problem hard to solve using basis expansion!

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Changing the resolution scale

 $\langle \psi_F(\lambda_0) | O(\lambda_0) | \psi_A(\lambda_0) \rangle = \left\langle \psi_F(\lambda_0) | U^{\dagger} U(\lambda) O(\lambda_0) U^{\dagger} U(\lambda) | \psi_A(\lambda) \right\rangle$ $= \left\langle \psi_F(\lambda) | O(\lambda) | \psi_A(\lambda) \right\rangle$

with

 $|\psi(\lambda)\rangle = U(\lambda) |\psi(\lambda_0)\rangle$ $O(\lambda) = U^{\dagger}(\lambda)O(\lambda_0)U(\lambda)$

 $U(\lambda)U^{\dagger}(\lambda) = 1$

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Convenient to choose resolution scale λ such that

- wave functions include only momentum scales that are constrained by scattering phase shifts (reduction of scheme dependence)
- nuclear structure calculations are simplified

• generate unitary transformation which decouples low- and high momenta:

 $H_{\lambda} = U_{\lambda} H U_{\lambda}^{\dagger}$ with the resolution parameter λ

$$\underbrace{\frac{dH_{\lambda}}{d\lambda} = [\eta_{\lambda}, H_{\lambda}]}$$

- generator η_{λ} can be chosen and tailored to different applications
- observables are preserved due to unitarity of transformation

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Systematic decoupling of high-momentum physics: the Similarity Renormalization Group



- elimination of coupling between low- and high momentum components,
 —> simplified many-body calculations!
- observables unaffected by resolution change (for exact calculations)
- residual resolution dependences can be used as tool to test calculations

Not the full story:

RG transformations also change three-body (and higher-body) interactions...















• detection of knocked out pairs with large relative momenta

• excess of np pairs over pp pairs

Subedi et al., Science 320, 1476 (2008)







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Explanation in terms of low-momentum interactions?





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Subedi et al., Science 320, 1476 (2008)

Explanation in terms of low-momentum interactions?

Vertex depends on the resolution! One-body current and SRC changes to two-body current and simple wave function.

Furnstahl, KH, Rep. Prog. Phys. 76, 126301 (2013)



 \Rightarrow quantities like momentum distributions are generally scale dependent:



- applies generally to all quantities like spectroscopic factors, short-range corr.,...
- consistency requires consistent RG evolution of reaction and structure parts
- key for all momenta involving high-momentum components

Evolved density operator in the deuteron

investigate $\langle \psi_D(\lambda) | U(\lambda) a_q^{\dagger} a_q U^{\dagger}(\lambda) | \psi_D(\lambda) \rangle$



Anderson, Bogner, Furnstahl, Perry, PRC 82, 054001 (2010)

- for low-momentum operators RG evolution provides only small corrections
- for high-momentum operators induced two-body contributions at small momenta completely dominate contribution at small resolution scales

Evolved density operator in the deuteron



- short-distance correlations in wave function at very resolution dependent
- perfect invariance of momentum distribution function with evolved density operator
- $U_{\lambda}(k,q)$ factorizes for $k < \lambda$ and $q \gg \lambda$: $U_{\lambda}(k,q) \approx K_{\lambda}(k)Q_{\lambda}(q)$

Electron inclusive cross sections at high momentum: Scaling in nuclear systems



idea: high-energy electrons probe the high-momentum part/small distance part of the nucleon wave function
the form of the high-momentum tail should be universal since it is dominated by short-range correlations
universality shows up as a plateaus in ratios of cross sections

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Egiyan et al. PRL 96, 1082501 (2006)

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the form of the high-momentum tail

should be universal since it is dominated by short-range correlations

• universality shows up as a plateaus in ratios of cross sections

• kinematically, 2-body correlations are limited to $x_B \sim 1$ to 2 and 3-body correlations to $x_B \sim 2$ to 3

Scaling in nuclear systems



- scaling behavior of momentum distribution function: $\rho_{NN}(q, Q = 0) \approx C_A \times \rho_{NN,Deuteron}(q, Q = 0)$ at large q
- dominance of np pairs over pp pairs
- "hard" (high resolution) interaction used
- dominance explained by short-range tensor forces

Nuclear scaling within chiral EFT

 $P \gg \Lambda_b \gg m_\pi \Rightarrow F_2^A(x, Q^2) \sim g_2(A, \Lambda) f_2(x, Q^2, \Lambda)$



Detailed comparison of experiment and theory

N²LO (
$$R_0 = 1.0 - 1.2 \text{ fm}$$
)AV18+UIXExp³H2.1(2) - 2.3(3)2.0(4)³He2.1(2) - 2.1(3)2.0(4)⁴He3.8(7) - 4.2(8)3.4(3)

Nuclear scaling at low resolution

 $\langle \psi_{\lambda} | O_{\lambda} | \psi_{\lambda} \rangle$ factorizes into a low-momentum structure and a **universal** high momentum part if the initial operator only weakly couples low and high momenta \longrightarrow explains scaling!

key:
$$U_{\lambda}(k, q) \approx K(k)Q(q)$$
 for $k < \lambda$ and $q \gg \lambda$
factorization!

That leads to:

$$\begin{aligned} \langle \psi_{\lambda} | O_{\lambda} | \psi_{\lambda} \rangle &= \int_{0}^{\lambda} dk \, dk' \int_{0}^{\infty} dq \, dq' \psi^{\dagger}(k) U_{\lambda}(k,q) O(q,q') U_{\lambda}(q',k') \cdot \psi_{\lambda}(k') \\ &\approx \int_{0}^{\lambda} dk \, dk' \, \psi_{\lambda}^{\dagger}(k) \left[\int_{0}^{\lambda} dq \, dq' K(k) K(q) O(q,q') K(q') K(k') + I_{QOQ} K(k) K(k') \right] \psi_{\lambda}(k') \end{aligned}$$

with the **universal** quantity:

$$I_{QOQ} = \int_{\lambda}^{\infty} dq \, dq' Q(q) O(q, q') Q(q')$$

valid if initial operator weakly couples low and high momenta





consider initial one-body current:

$$\begin{aligned} \langle \mathbf{k}_1 \, T_1 | \, J_0(\mathbf{q}) \, | \, \mathbf{k}_2 \, T &= 0 \rangle \\ &= \frac{1}{2} \left(G_E^p + (-1)^{T_1} G_E^n \right) \delta(\mathbf{k}_1 - \mathbf{k}_2 - \mathbf{q}/2) \\ &+ \frac{1}{2} \left((-1)^{T_1} G_E^p + G_E^n \right) \delta(\mathbf{k}_1 - \mathbf{k}_2 + \mathbf{q}/2) \end{aligned}$$

$$\langle \psi_f | J_0 | \psi_i \rangle = \underbrace{\langle \phi | J_0 | \psi_i \rangle}_{\text{IA}} + \underbrace{\langle \phi | t^{\dagger} G_0^{\dagger} J_0 | \psi_i \rangle}_{\text{FSI}}$$

$$= \langle \psi_f | U^{\dagger} U J_0 U^{\dagger} U | \psi_i \rangle$$

$$= \langle \psi_f^{\lambda} | J_0^{\lambda} | \psi_i^{\lambda} \rangle$$

$$\mathcal{T}_{S,m_{s_f},\mu,m_{J_d}} = -\pi \sqrt{2\alpha |\mathbf{p}'| E_p E_d / M_d} \langle \psi_f | J_\mu(\mathbf{q}) | \psi_i \rangle$$

study longitudinal structure function

$$f_{L} = \sum_{\substack{S_{f}, m_{s_{f}} \\ m_{J_{d}}}} \mathcal{T}_{S_{f}, m_{s_{f}}, \mu = 0, m_{J_{d}}}(\theta', \varphi') \, \mathcal{T}^{*}_{S_{f}, m_{s_{f}}, \mu = 0, m_{J_{d}}}(\theta', \varphi')$$





Summary

- for ab-initio studies of reactions it is crucial to treat structure and reaction part consistently and simultaneously
- theoretical interpretation sensitively depends on the resolution scale
- resolution scale change will shift contributions between structure and reaction parts
- deep inelastic cross sections usually explained in terms of short-range correlations, scheme dependent, all observables can also be explained by separation of scales and factorization
- studied deuteron disintegration based on RG evolved interactions and currents
 - + found perfect RG invariance of longitudinal structure function
 - + impact of RG evolution strongly depends on kinematics