

Darmstadt Lecture 2 – Equation of State

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Pressure - Radius Correlation

Newtonian polytrope: $p = K\rho^{1+1/n}$

$$\Rightarrow R \propto K^{n/(3-n)} M^{(1-n)/(3-n)}$$

Realistic EOS: $n \simeq 1 \Rightarrow R \propto K^{1/2} M^0$

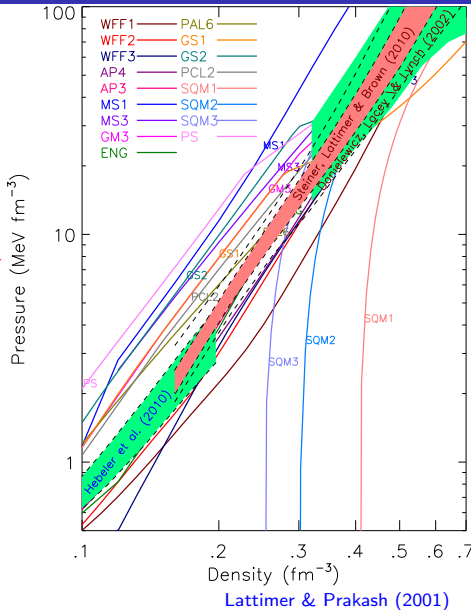
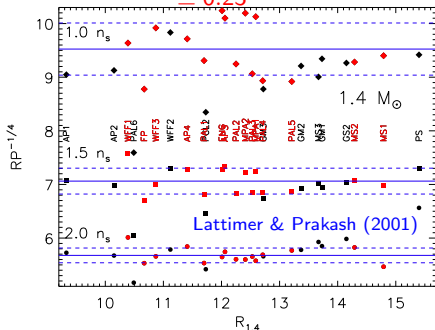
GR phenomenological result:

$$R \propto K^{1/4} \propto p_f^{1/4} \rho_f^{-1/2}$$

Buchdahl motivation:

$$2R^2 p_*(1 - 2\beta) = \pi(1 - \beta)^2 \Rightarrow$$

$$\left. \frac{d \ln R}{d \ln \rho} \right|_{n,M} = \frac{(1 - \beta)(1 - 2\beta)(\sqrt{p_*} - 10\sqrt{p})}{2(1 - 3\beta + 3\beta^2)(\sqrt{p_*} + 2\sqrt{p})} \approx 0.23$$



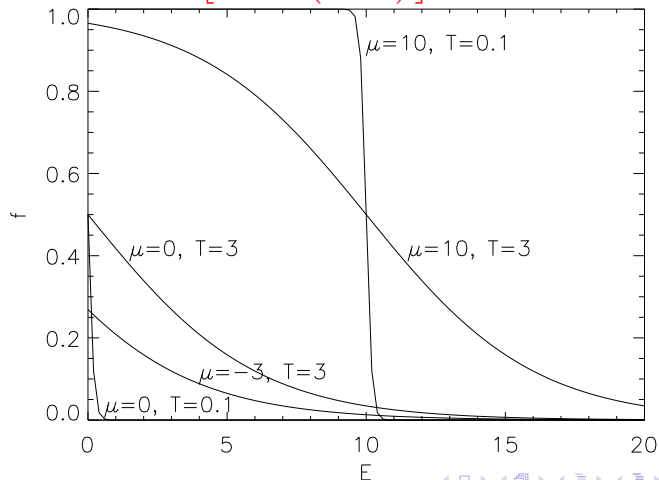
EOS of an Ideal Fermion Gas

Single particle energy

$$E^2 = m^2 c^4 + p^2 c^2$$

Occupation index

$$f = \left[1 + \exp\left(\frac{E - \mu}{T}\right) \right]^{-1}$$



Thermodynamics

Number and energy densities

$$n = \frac{g}{h^3} \int_0^\infty f d^3 p; \quad \varepsilon = \frac{g}{h^3} \int_0^\infty E f d^3 p; \quad \mu = \left(\frac{\partial F}{\partial n} \right)_T$$

Landau quasi-particle entropy formula:

$$ns = -\frac{g}{h^3} \int_0^\infty [f \ln f + (1-f) \ln(1-f)] d^3 p$$

Pressure

$$P = \frac{g}{3h^3} \int_0^\infty p \frac{\partial E}{\partial p} f d^3 p = n^2 \left(\frac{\partial(F/n)}{\partial n} \right)_T = \mu n - \varepsilon + nTs = \mu n - F$$

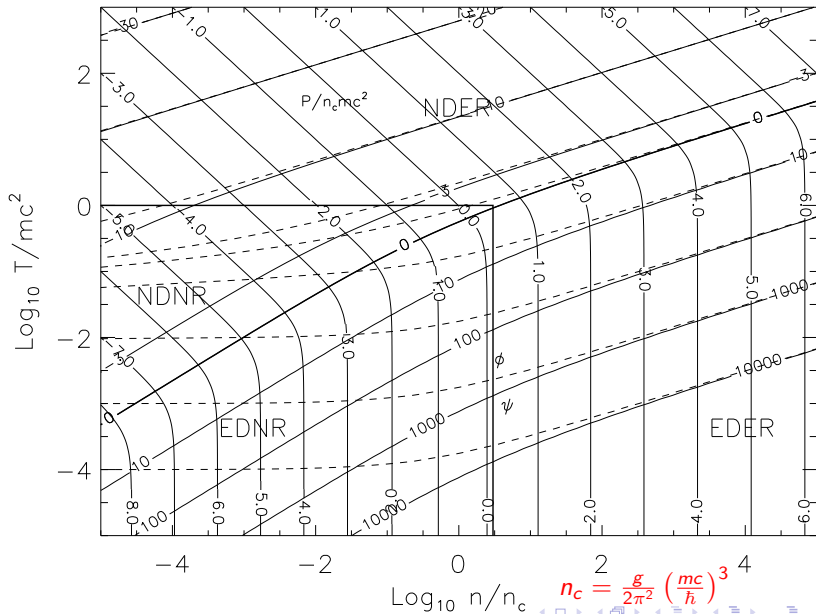
$$n = \left(\frac{\partial P}{\partial \mu} \right)_T; \quad ns = \left(\frac{\partial P}{\partial T} \right)_\mu$$

Degeneracy parameters

$$\phi = \frac{\mu}{T} = \psi + \frac{mc^2}{T}; \quad \left(\frac{\partial P}{\partial T} \right)_\phi = n(s + \phi); \quad \left(\frac{\partial P}{\partial \psi} \right)_T = n(s + \psi)$$

Pressure and Degeneracy Parameters

Fermions



Non-relativistic $p \ll mc$

$$x = p^2/(2mT), \quad \psi = (\mu - mc^2)/T$$

$$n = \frac{g(2mT)^{3/2}}{4\pi^2\hbar^3} \int_0^\infty \frac{x^{1/2}}{1 + e^{x-\psi}} dx = \frac{g(2mT)^{3/2}}{4\pi^2\hbar^3} F_{1/2}(\psi)$$

$$\varepsilon = nmc^2 + \frac{gT(2mT)^{3/2}}{4\pi^2\hbar^3} F_{3/2}(\psi)$$

$$P = \frac{2}{3}(\varepsilon - nmc^2), \quad s = \frac{5F_{3/2}(\psi)}{3F_{1/2}(\psi)} - \psi$$

Relativistic $p \gg mc$

$$n = \frac{g}{2\pi^2} \left(\frac{T}{\hbar c} \right)^3 F_2(\phi)$$

$$\varepsilon = \frac{gT}{2\pi^2} \left(\frac{T}{\hbar c} \right)^3 F_3(\phi)$$

$$P = \frac{1}{3}\varepsilon, \quad s = \frac{4F_3(\phi)}{3F_2(\phi)} - \phi$$

$$\left\{ \begin{array}{l} n = \frac{g}{6\pi^2} \left(\frac{\mu}{\hbar c}\right)^3 \left[1 + \left(\frac{\pi}{\phi}\right)^2 + \dots\right], \quad s = \frac{\pi^2}{\phi} + \dots \\ P = \frac{1}{3}\varepsilon = \frac{1}{4}n\mu \left[1 + \left(\frac{\pi}{\phi}\right)^2 + \dots\right] \end{array} \right\} \text{EDER}$$

$$\left\{ \begin{array}{l} n = \frac{g}{6\pi^2} \left(\frac{2m\psi T}{\hbar^2}\right)^{3/2} \left[1 + \frac{1}{8} \left(\frac{\pi}{\psi}\right)^2 + \dots\right], \quad s = \frac{\pi^2}{2\psi} + \dots \\ P = \frac{2}{3}(\varepsilon - nmc^2) = \frac{2n\psi T}{5} \left[1 + \frac{1}{2} \left(\frac{\pi}{\psi}\right)^2 + \dots\right] \end{array} \right\} \text{EDNR}$$

$$\left\{ \begin{array}{l} n = \frac{g}{\pi^2} \left(\frac{T}{\hbar c}\right)^3 e^{\phi}, \quad s = 4 - \phi \\ P = \frac{1}{3}\varepsilon = nT \end{array} \right\} \text{NDER}$$

$$\left\{ \begin{array}{l} n = g \left(\frac{mT}{2\pi\hbar^2}\right)^{3/2} e^{\psi}; \quad s = \frac{5}{2} - \psi \\ P = \frac{2}{3}(\varepsilon - nmc^2) = nT \end{array} \right\} \text{NDNR}$$

Gas With Pairs in Extremely Relativistic Limit

$$\mu = \mu^+ = -\mu^-,$$

$$n = n^+ - n^- = \frac{g}{6\pi^2} \left(\frac{\mu}{\hbar c}\right)^3 \left[1 + \left(\frac{\pi T}{\mu}\right)^2\right],$$

$$P = P^+ + P^- = \frac{\varepsilon}{3} = \frac{g\mu}{24\pi^2} \left(\frac{\mu}{\hbar c}\right)^3 \left[1 + 2\left(\frac{\pi T}{\mu}\right)^2 + \frac{7}{15}\left(\frac{\pi T}{\mu}\right)^4\right],$$

$$s = s^+ + s^- = \frac{gT}{6n} \left(\frac{\mu}{\hbar c}\right)^2 \left[1 + \frac{7}{15}\left(\frac{\pi T}{\mu}\right)^2\right].$$

Cubic Solution : $\mu(n, T)$

$$\mu = r - q/r, \quad r = \left[(q^3 + t^2)^{1/2} + t\right]^{1/3},$$

$$t = \frac{3\pi^2}{g} n(\hbar c)^3, \quad q = \frac{(\pi T)^2}{3}.$$

Interacting Non-Relativistic Fermi Gases

Single particle energy (non-relativistic), $t = n, p$

$$E_t = m_t c^2 + \frac{p^2}{2m_t^*(n_n, n_p)} + V_t(n_n, n_p)$$

$$\frac{\hbar^2}{2m_t^*} = \frac{\delta \varepsilon}{\delta \tau_t}, \quad V_t = \frac{\delta \varepsilon}{\delta n_t}, \quad f_t = \left[1 + \exp\left(\frac{E_t - \mu_t}{T}\right) \right]^{-1}$$

Number and kinetic densities

$$n_t = \frac{1}{2\pi^2 \hbar^2} \left(\frac{2m_t^* T}{\hbar^2} \right)^{3/2} F_{1/2}(\eta_t); \quad \tau_t = \frac{1}{2\pi^2 \hbar^2} \left(\frac{2m_t^* T}{\hbar^2} \right)^{5/2} F_{3/2}(\eta_t)$$

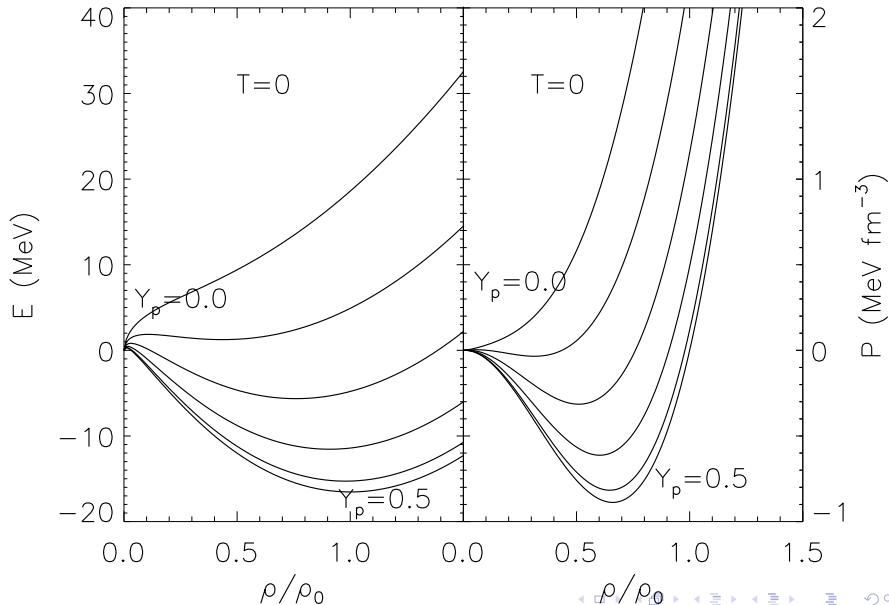
$$\eta_t = \frac{\mu_t - m_t c^2 - V_t}{T}$$

Thermodynamic quantities

$$\varepsilon = \sum_t \frac{\hbar^2}{2m_t^*(n_n, n_p)} \tau_t + U(n_n, n_p); \quad V_t = \frac{\hbar^2}{2} \sum_s \tau_s \frac{\partial (m_s^*)^{-1}}{\partial n_s} + \frac{\partial U}{\partial n_t}$$

$$P = \sum_t \left[n_t V_t + \frac{\hbar^2 \tau_t}{3m_t^*} \right] - U, \quad ns = \sum_t \left[\frac{5\hbar^2 \tau_t}{6m_t^* T} - n_t \eta_t \right]$$

Bulk Matter Energy and Pressure



Schematic Free Energy Density

$F(n, x, T)$; n : number density; x : proton fraction; T : temperature

$n_s \simeq 0.16 \pm 0.01 \text{ fm}^{-3}$: nuclear saturation density

$B \simeq 16 \pm 1 \text{ MeV}$: saturation binding energy

$K_s \simeq 240 \pm 20 \text{ MeV}$: incompressibility

$S_v \simeq 30 \pm 6 \text{ MeV}$: bulk symmetry energy

$L \simeq 60 \pm 60 \text{ MeV}$: symmetry stiffness

$a \simeq 0.065 \pm 0.010 \text{ MeV}^{-1}$: bulk level density parameter

$K'_s \simeq -200 \pm 200 \text{ MeV}$:
skewness

$K_{sym} \simeq -300 \pm 300 \text{ MeV}$:
symmetry incompressibility

$$\begin{aligned} F &= n \left[-B + \frac{K}{18} \left(1 - \frac{n}{n_s} \right)^2 + S_v \frac{n}{n_s} (1 - 2x)^2 - a \left(\frac{n_s}{n} \right)^{2/3} T^2 \right] \\ P &= n^2 \frac{\partial(F/n)}{\partial n} = \frac{n^2}{n_s} \left[\frac{K}{9} \left(\frac{n}{n_s} - 1 \right) + S_v (1 - 2x)^2 \right] + \frac{2an}{3} \left(\frac{n_s}{n} \right)^{2/3} T^2 \\ \mu_n &= \frac{\partial F}{\partial n} - \frac{x}{n} \frac{\partial F}{\partial x} \\ &= -B + \frac{K}{18} \left(1 - \frac{n}{n_s} \right) \left(1 - 3 \frac{n}{n_s} \right) + 2S_v \frac{n}{n_s} (1 - 2x) - \frac{a}{3} \left(\frac{n_s}{n} \right)^{2/3} T^2 \\ \hat{\mu} &= -\frac{1}{n} \frac{\partial F}{\partial x} = \mu_n - \mu_p = 4S_v \frac{n}{n_s} (1 - 2x) \\ s &= -\frac{1}{n} \frac{\partial F}{\partial T} = 2a \left(\frac{n_s}{n} \right)^{2/3} T; \quad \varepsilon = F + nTs \end{aligned}$$

Phase Coexistence

Negative pressure: matter is unstable to separating into two phases of different densities (and possibly proton fractions). Physically, this represents coexistence of nuclei and vapor. Neglecting finite-size effects, bulk coexistence approximates the EOS at subnuclear densities.

Free Energy Minimization With Two Phases

$$\begin{aligned}F &= \epsilon - nTs = uF_I + (1-u)F_{II}, \\n &= un_I + (1-u)n_{II}, \\nY_e &= ux_I n_I + (1-u)x_{II} n_{II}.\end{aligned}$$

$$n_{II} = \frac{n - un_I}{1-u}, \quad x_{II} = \frac{nY_e - un_I x_I}{n - un_I}$$

$$\frac{dF}{dn_I} = u \frac{\partial F_I}{\partial n_I} + (1-u) \frac{\partial F_{II}}{\partial n_{II}} \left(\frac{-u}{1-u} \right) = u(\mu_{n,I} - \mu_{n,II})$$

$$\frac{dF}{dx_I} = u \frac{\partial F_I}{\partial x_I} + (1-u) \frac{\partial F_{II}}{\partial x_{II}} \left(\frac{-un_I}{n - un_I} \right) = un_I(\hat{\mu}_I - \hat{\mu}_{II})$$

$$\frac{dF}{du} = F_I - F_{II} + (1-u) \left[\frac{\partial F_{II}}{\partial n_{II}} \left(\frac{n_{II} - n_I}{1-u} \right) + \frac{\partial F_{II}}{\partial x_{II}} \left(\frac{-un_I}{n - un_I} \right) \right]$$

$$\implies \mu_{nI} = \mu_{nII}, \quad \mu_{pI} = \mu_{pII}, \quad P_I = P_{II}$$

Critical Point
($Y_e = 0.5$)

$$\left(\frac{\partial P}{\partial n}\right)_T = \left(\frac{\partial^2 P}{\partial n^2}\right)_T = 0$$

$$n_c = \frac{5}{12} n_s$$

$$T_c = \left(\frac{5}{12}\right)^{1/3} \left(\frac{5K}{32a}\right)^{1/2}$$

$$s_c = \left(\frac{12}{5}\right)^{1/3} \left(\frac{5Ka}{8}\right)^{1/2}$$

