

Darmstadt Lecture 7 – Proto-Neutron Stars

James Lattimer

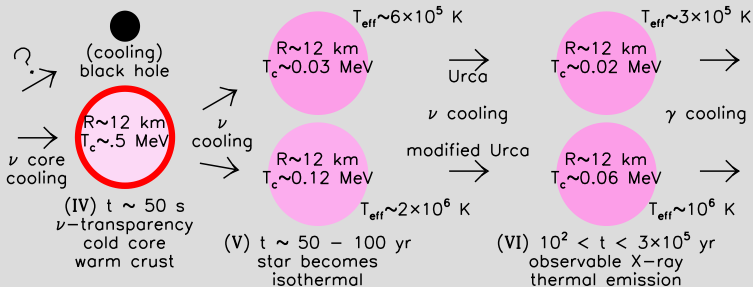
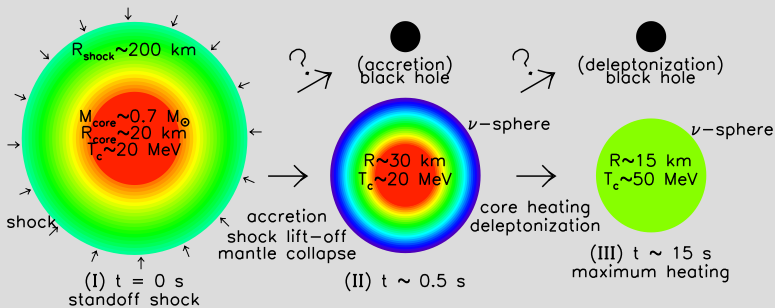
Department of Physics & Astronomy
449 ESS Bldg.
Stony Brook University

July 16, 2019

Darmstadt Lecture 7 – Proto-Neutron Stars

James.Lattimer@Stonybrook.edu

Proto-Neutron Stars



Proto-Neutron Star Evolution

Consider the Newtonian case. The electron neutrino number flux F_ν and ν energy flux L_ν change the lepton number Y_L and entropy s :

$$\begin{aligned}n \frac{dY_L}{dt} &= n \frac{dY_e}{dt} + \frac{d(Y_\nu - Y_{\bar{\nu}})}{dt} = \frac{1}{r^2} \frac{\partial}{\partial r} r^2 F_\nu \\nT \frac{ds}{dt} &= -\frac{1}{4\pi r^2} \frac{\partial(L_\nu + L_{\bar{\nu}})}{\partial r} - n \sum_{n,p,e,\nu} \mu_i \frac{dY_i}{dt}.\end{aligned}$$

In the diffusion approximation, fluxes are driven by density gradients:

$$\begin{aligned}F_\nu &= -\int_0^\infty \frac{c}{3} \left(\lambda_N \frac{\partial n_\nu(E)}{\partial r} - \bar{\lambda}_N \frac{\partial n_{\bar{\nu}}(E)}{\partial r} \right) dE, \\L_\nu &= -\int_0^\infty 4\pi r^2 \sum_i^6 \frac{c \lambda_E^i}{3} \frac{\partial \epsilon_i(E)}{\partial r} dE.\end{aligned}$$

λ_N and λ_E^i 's are mean free paths for number and energy transport, respectively. $n_\nu(E)$ and $\epsilon_i(E)$ are the ν_e number and energy densities of species $i = e, \mu, \tau$ and their antiparticles at neutrino energy E .

There are two main sources of opacity:

ν -nucleon absorption, affects only $\nu_e, \bar{\nu}_e$.

$\lambda_\nu \simeq \bar{\lambda}_\nu \simeq \lambda_0 (E_0/E)^2$; $\lambda_0 \approx 5$ cm; $E_0 \approx 1$ MeV

$\nu - e$ scattering, affects all flavors. $\lambda_E^i \simeq \lambda_1^i (E_0/E)^2$; $\lambda_1 \approx 4$ km

Proto-Neutron Stars – Analytic Analysis

Neutrino fluid

$$n_\nu(E) = \frac{E^2}{2\pi^3(\hbar c)^3} f_\nu(E), \quad f_\nu(E) = \left[1 + e^{(E-\mu_\nu)/T}\right]^{-1}$$

$$n(Y_\nu - Y_{\bar{\nu}}) = \frac{\mu_\nu^3 + \pi^2 T^2 \mu}{6\pi^2(\hbar c)^3}$$

$$\epsilon_\nu(E) = E(n_\nu(E) + n_{\bar{\nu}}(E))$$

Diffusion approximation

$$F_\nu = -\frac{c}{3} \int_0^\infty \left[\lambda_N \frac{\partial n_\nu(E)}{\partial r} - \bar{\lambda}_N \frac{\partial n_{\bar{\nu}}(E)}{\partial r} \right] dE$$
$$\simeq -\frac{c\lambda_0 E_0^2 T}{6\pi^2(\hbar c)^3} \frac{\partial T[F_0(\mu_\nu/T) - F_0(-\mu_\nu/T)]}{\partial r} = -\frac{c\lambda_0 E_0^2}{6\pi^2(\hbar c)^3} \frac{\partial \mu_\nu}{\partial r},$$

$$L_\nu = -4\pi r^2 \int_0^\infty \sum_i^6 \frac{c\lambda_E^i}{3} \frac{\partial \epsilon_i(E)}{\partial r} dE$$
$$\simeq -\frac{r^2 c}{3\pi(\hbar c)^3} \sum_i^3 \lambda_1^i E_0^{i2} \frac{\partial T^2[F_1(\mu_\nu/T) + F_1(-\mu_\nu/T)]}{\partial r}$$
$$= -\frac{r^2 c}{3\pi(\hbar c)^3} \sum_i^3 \lambda_1^i E_0^{i2} \frac{\partial}{\partial r} \left(\mu_\nu^2 + \frac{\pi^2 T^2}{3} \right).$$

The Deleptonization of a Proto-Neutron Star

Energy transport dominated by degenerate electron neutrinos propagating through degenerate matter.

Number transport equation

$$\begin{aligned}\frac{\partial n Y_L}{\partial t} &= \frac{\partial Y_L}{\partial Y_\nu} \frac{1}{6\pi^2(\hbar c)^3} \frac{\partial \mu_\nu^3}{\partial t} = \frac{1}{r^2} \frac{\partial}{\partial r} r^2 F_\nu \\ &= -\frac{c\lambda_0 E_0^2}{6\pi^2(\hbar c)^3} \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial \mu_\nu}{\partial r} \right)\end{aligned}$$

$$\frac{\partial Y_L}{\partial Y_\nu} \simeq 5, \quad \mu_\nu = \mu_{\nu,0} \psi(x) \phi(t), \quad x = x_1 r/R$$

$$\tau_D \frac{d\phi^2}{dt} = -\frac{1}{x^2 \psi^3} \frac{\partial}{\partial x} \left[x^2 \frac{\partial \psi}{\partial x} \right] = -1$$

Spatial solution is Lane-Emden function of index 3, ψ_3 .

$$\phi = \sqrt{1 - \frac{t}{\tau_D}}, \quad \psi_3(x_1) = 0 \implies x_1 = 6.897$$

$$\tau_D = \frac{6}{c\lambda_0} \frac{\partial Y_L}{\partial Y_\nu} \left(\frac{\mu_{\nu,0}}{E_0} \right)^2 \left(\frac{R}{x_1} \right)^2 \approx 5 - 10 \text{ s}$$

Heating During Deleptonization

Beta equilibrium:

$$\sum_j \mu_j dY_j = (-\mu_n + \mu_p + \mu_e - \mu_\nu) dY_e + \mu_\nu dY_L = \mu_\nu dY_L$$

Energy transport equation:

$$nT \frac{\partial s}{\partial t} = -\frac{1}{4\pi r^2} \frac{\partial L_\nu}{\partial r} - n \sum_j \mu_j \frac{\partial Y_j}{\partial t} \simeq -n\mu_\nu \frac{\partial Y_L}{\partial t}$$

$$s \simeq aT, \quad a \simeq 0.06 \text{ MeV}^{-1}$$

$$\frac{s}{a} \frac{ds}{dt} \simeq -\mu_\nu \frac{\partial Y_L}{\partial t} \simeq -3 \frac{\partial Y_L}{\partial Y_\nu} Y_{\nu,0} \mu_{\nu,0} \left(\frac{\mu_\nu}{\mu_{\nu,0}} \right)^3 \frac{\partial \mu_\nu / \mu_{\nu,0}}{\partial t}$$

$$s_f^2 - s_i^2 \simeq \frac{3a}{2} \left(\frac{\partial Y_L}{\partial Y_\nu} \right) \mu_{\nu,0} Y_{\nu,0}$$

$$s_i \sim 1, \quad s_f \sim 2.5, \quad T_f \simeq 50 \text{ MeV}$$

Following deleptonization, $\mu_\nu \ll T$

$$\begin{aligned}L_\nu &= -4\pi r^2 \frac{c}{6\pi^2(\hbar c)^3} \frac{\partial}{\partial r} \int_0^\infty \sum_i^6 \lambda_E^i E_0^{i2} E f_\nu dE \\ &= -4\pi r^2 \frac{c\lambda_1 E_0^2}{\pi^2(\hbar c)^3} \frac{\partial}{\partial r} \frac{\pi^2 T^2}{3}.\end{aligned}$$

Energy transport equation:

$$\begin{aligned}nT \frac{\partial s}{\partial t} &= -\frac{1}{4\pi r^2} \frac{\partial L_\nu}{\partial r} - n \sum_j \mu_j \frac{dY_j}{dt} \simeq -\frac{1}{4\pi r^2} \frac{\partial L_\nu}{\partial r} \\ naT \frac{\partial T}{\partial t} &= \frac{c\lambda_1 E_0^2}{3(\hbar c)^3} \frac{1}{r^2} \frac{\partial}{\partial r} r^2 \frac{\partial T^2}{\partial r}.\end{aligned}$$

Assume $T = T_0\phi(t)\psi(x)$, $T_0 \sim 50$ MeV.

$$\frac{\tau_c}{\phi} \frac{d\phi}{dt} = -1 = \frac{1}{\psi^2 x^2} \frac{\partial}{\partial x} x^2 \frac{\partial \psi^2}{\partial x}.$$

ψ^2 is Lane-Emden function of index 1: $\psi^2 = \psi_1 = \sin x/x$

$$\psi_1(x_1) = 0 \implies x_1 = \pi, \quad \left(x \frac{\partial \psi_1}{\partial x} \right)_{x_1} = -1.$$

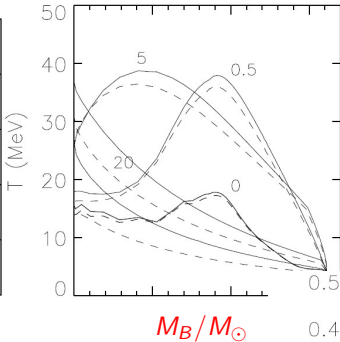
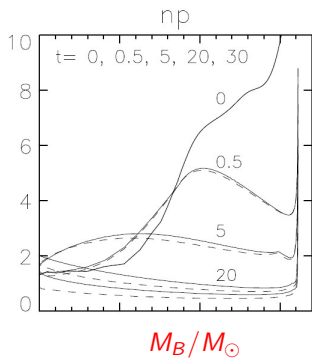
$$\phi = \exp\left(\frac{t_0 - t}{\tau_c}\right), \quad \tau_c = \frac{3(\hbar c)^3 a n}{c \lambda_1 E_0^2} \left(\frac{R}{x_1}\right)^2 \approx 10 - 30 \text{ s}$$

Emergent Luminosity:

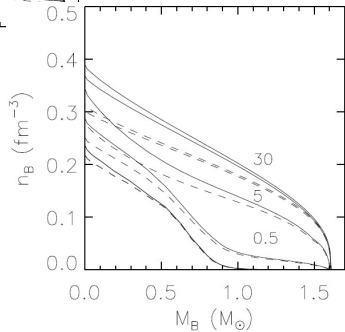
$$\begin{aligned} L_\nu(R, t) &= -4\pi R \frac{c \lambda_1 E_0^2 T_0^2}{3(\hbar c)^3} \left(x \frac{\partial \psi^2}{\partial x} \right)_{x_1} \phi^2(t) = \frac{c F_3(0)}{2(\hbar c)^3} R^2 T_{\nu, \text{eff}}(t)^4 \\ &= L_0 \phi(t)^2 \approx 1.2 \cdot 10^{51} \phi(t)^2 \text{ erg s}^{-1} \end{aligned}$$

$$T_{\nu, \text{eff}}(t) = \left(\frac{2L_0(\hbar c)^3}{c F_3(0) R^2} \right)^{1/4} \sqrt{\phi(t)} \simeq 4 \sqrt{\phi(t)} \text{ MeV}, \quad \langle E_\nu \rangle \sim 3 T_\nu.$$

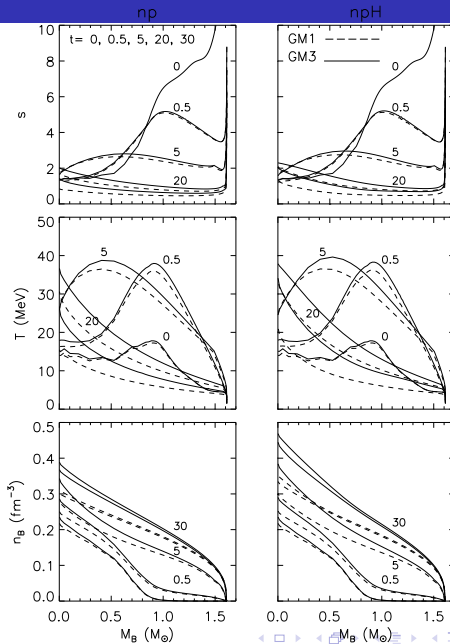
Model Simulations



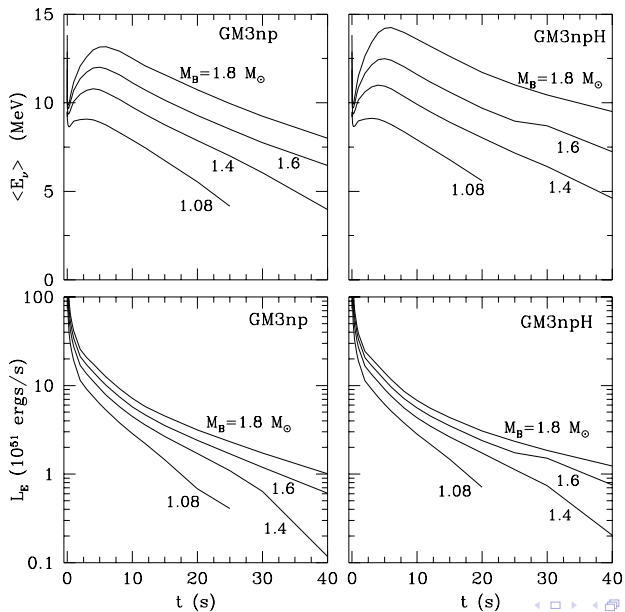
Pons et al.



Model Simulations



Model Simulations



Model Signal

